

## Quantum theory of light

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Sheet 6

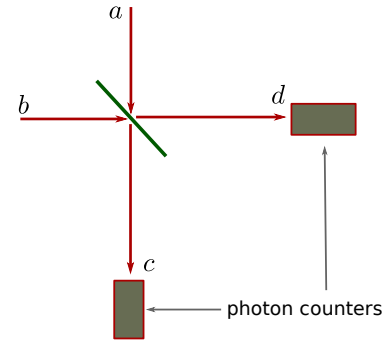
*Hand in:* Tue 03.12.2019 (questions marked as \* are optional)*Discussion date:* Mon 09.12.2019

### 14. Demonstrating squeezing by measuring the quadratures

In this problem we describe how the squeezed states are detected by means of the *balanced homodyne* setup. In this setup, the squeezed mode  $a$  is mixed with a mode  $b$  by a beam-splitter (see figure beside). The beam-splitter implements the transformation

$$c = \frac{1}{\sqrt{2}}(a + ib), d = \frac{1}{\sqrt{2}}(b + ia). \quad (1)$$

Suppose that all the modes are of the same frequency  $\omega$  (they are all derived from the same laser in practice.) The modes  $c$  and  $d$  are the subjected to photon counters which measure the photon numbers  $n_c$  and  $n_d$  and denote  $n_{cd} := n_c - n_d$ .



(a) [5pts] Show that  $n_{cd} = i(a^\dagger b - ab^\dagger)$ .

(b) [5pts] Suppose mode  $b$  is in a large coherent state  $|\beta\rangle$  with  $\beta = |\beta|e^{i\theta}$ , show that  $n_{cd}$  implements the measurement of a quadrature of mode  $a$ . The fluctuation thus allows one to determine the squeezed quadratures of the state.

### 15. Two-mode squeezing

In similarity to the one-mode squeezing operator, one can introduce the two-mode squeezing operator

$$S_2(\xi) = \exp(\xi^* ab - \xi a^\dagger b^\dagger) \quad (2)$$

with  $\xi = r e^{i\theta}$ , where  $a$  and  $b$  are two different modes. The two-mode squeezed vacuum is given by  $|\xi\rangle_2 = S_2(\xi) |0, 0\rangle$ .

(a) [5pts] Show that

$$S_2^\dagger(\xi) a S_2(\xi) = a \cosh r - e^{i\theta} b^\dagger \sinh r, \quad (3)$$

$$S_2^\dagger(\xi) b S_2(\xi) = b \cosh r - e^{i\theta} a^\dagger \sinh r. \quad (4)$$

(b) [10pts] Define the quadratures involving two modes  $X_1 = 1/\sqrt{8}(a + a^\dagger + b + b^\dagger)$ ,  $X_2 = 1/(\sqrt{8}i)(a - a^\dagger + b - b^\dagger)$ . Show that  $[X_1, X_2] = i/2$ . And show that for the two-mode squeezed vacuum  $|\xi\rangle_2$  with  $\theta = 0$ ,  $\langle X_1 \rangle = \langle X_2 \rangle = 0$  and  $\langle X_1^2 \rangle = e^{-2r}/4$ ,  $\langle X_2^2 \rangle = e^{2r}/4$ .

(c) [\* , 10pts] Show that the two-mode squeezed vacuum can be expanded in the number states as

$$|\xi\rangle_2 = \frac{1}{\cosh r} \sum_{n=0}^{+\infty} (-1)^n e^{in\theta} (\tanh r)^n |n, n\rangle. \quad (5)$$

Then show that the photon statistics of each mode is that of the thermal state with effective temperature  $T_{\text{eff}} = \hbar\omega_i / (2k_B \ln \cosh r)$ , where  $\omega_i$  are the frequencies of the modes.

- (c) [5pts] In the lecture, you learn how to produce a (one-mode) squeezed state by means of the *degenerate parametric down-conversion*. Following the same steps, describe how the two-mode squeezed vacuum can be generated by the *non-degenerate parametric down-conversion* with the Hamiltonian

$$H = \hbar\omega_a a^\dagger a + \hbar\omega_b b^\dagger b + \hbar\omega_p c^\dagger c + i\chi^{(2)}(abc^\dagger - a^\dagger b^\dagger c), \quad (6)$$

where  $c$  denotes the pumping mode, and  $\chi^{(2)}$  is the coupling.