

Reciprocal Space and x-ray diffraction

Due to 3D the periodic arrangement of atoms (molecules) in a crystal one can find “dense packed planes” = **lattice planes**, in certain directions of space. Assuming one of these lattice planes makes an intersection point with the a, b, c axes with the ratios ma;nb;pc defining the direction of the normal of this plane by $\vec{r} = ma + nb + pc$ the inverse ratio normalized to integer values $\frac{1}{m}; \frac{1}{n}; \frac{1}{p} = h, k, l$ define the **Miller indices** h, k, l which are integer multiples of

a vector $\mathbf{k}$ in reciprocal space $\vec{k} = h\mathbf{a}^* + k\mathbf{b}^* + l\mathbf{c}^*$. Miller indices of the main plane cutting the a-axis of a cube are (100), its normal vector is [100]. The shortest distance between two neighboured lattice planes is the **lattice plane distance d**; d is a function of h,k,l and the lattice parameters. In cubic system it is given by $d = \frac{a}{\sqrt{h^2 + k^2 + l^2}}$; in hexagonal system it is

$$d = \frac{a}{\sqrt{\frac{4}{3}(h^2 + k^2 + hk) + \frac{a^2}{c^2}l^2}};$$

The **reciprocal lattice** is related to the periodicity of the crystal lattice and can be considered as it's Fourier-transformation. The basis vectors $\mathbf{b}_i$ of the reciprocal lattice are related to the basis vectors $\mathbf{a}_i$ by the relation

$$\mathbf{b}_i = 2\pi \varepsilon_{ijk} \frac{\mathbf{a}_j \times \mathbf{a}_k}{\mathbf{a}_1 \cdot \mathbf{a}_2 \times \mathbf{a}_3} \text{ and fulfill the relation } \mathbf{a}_i \cdot \mathbf{b}_j = 2\pi \delta_{ij}.$$

In crystal systems with 90° angles between axes $\mathbf{b}_i^*$ is parallel to $\mathbf{a}_i$ and reciprocal in length $b_i^* = 2\pi/a_i$. Because the wave length of x-rays are in same order of magnitude as lattice

spacing $\mathbf{a}_i$ the arrangement of atoms in a crystal can be studied by an x-ray diffraction experiment. The interference condition of a x-ray wave with a 3D crystal has been derived by Max v. Laue in 1912. If $\alpha_0, \beta_0, \gamma_0$ are the angles of incidence and α, β, γ the angles of exit with respect to the axes a,b,c one can find 3 conditions to be fulfilled simultaneously in order to obtain a diffraction maximum in 3D space: Maxima appear whenever:

$g_x (\cos \alpha - \cos \alpha_0) = \pm h \lambda$; $g_y (\cos \beta - \cos \beta_0) = \pm k \lambda$; $g_z (\cos \gamma - \cos \gamma_0) = \pm l \lambda$
For a certain set of Miller indices h, k, l this is possible only for a single wave length.

$$\lambda = \frac{2 \frac{N_z}{g_z}}{\left(\frac{N_x}{g_x}\right)^2 + \left(\frac{N_y}{g_y}\right)^2 + \left(\frac{N_z}{g_z}\right)^2}$$

Laue's interference condition

An alternative relation to obtain an interference maximum has been derived by Bragg:

$n\lambda = 2d \sin \Theta$, where Θ is the reflection angle with respect to a set of lattice planes with spacing d. A graphic tool to solve Laue's and Bragg's equations is the concept of Ewald's sphere.

The Intensities of the Bragg reflections (see Laue Interference) depend on the number of unit cells within the crystal (**shape factor, S**)

$$S = \frac{\sin^2(N \frac{\pi d}{\lambda} \sin \alpha)}{\sin^2(\frac{\pi d}{\lambda} \sin \alpha)^2}$$

and the relative coordinates of atoms within the unit cell **structure factors** $F = \sum_n f_n \exp(-2\pi i(hx_i + ky_i + lz_i))$. The scattering

intensity is proportional to S^2 and F^2 . Typical diffraction experiments executed in Siegen's Lab. practicum) are the Laue diffraction and Debye-Scherrer-Experiment. The Solid State Physics representation of a primitive reciprocal unit cell is the **Brillion zone**.