

# Quantum gates in rare-earth-ion doped crystals



Atia Amari, Brian Julsgaard  
Stefan Kröll, Lars Rippe  
Andreas Walther, Yan Ying

**Knut och Alice  
Wallenbergs  
Stiftelse**



# Outline

- Rare-earth-ion doped crystals as quantum computer hardware
- Experimental results
- Current status and outlook

# Requirements for quantum computing

- **Coherent two-level systems acting as qubits**
- Possibility to manipulate the qubits individually (single qubit operations)
- Coupling between any two qubits (two-bit gates)
- Possibility for reliable read-out of the individual qubits
- Scalability

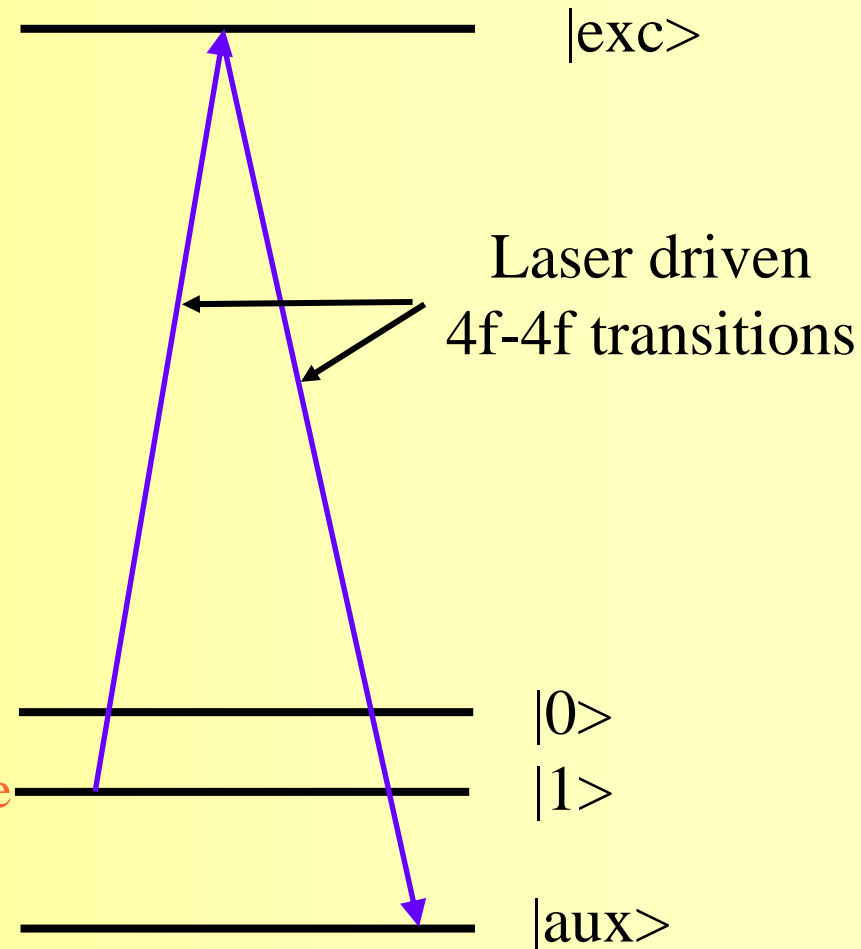
# The rare-earth-ions hyperfine states are used as qubit states

- Long coherence times of the optical transitions (100  $\mu\text{s}$  to 6 ms)

- At 4 Kelvin the ground state hyperfine levels have ms-s coherence times and very long lifetimes (~ hours)

- The duration of a laser  $\pi$ -pulse,  $\sim 1 \mu\text{s}$

Ground state  
with hyperfine  
splitting

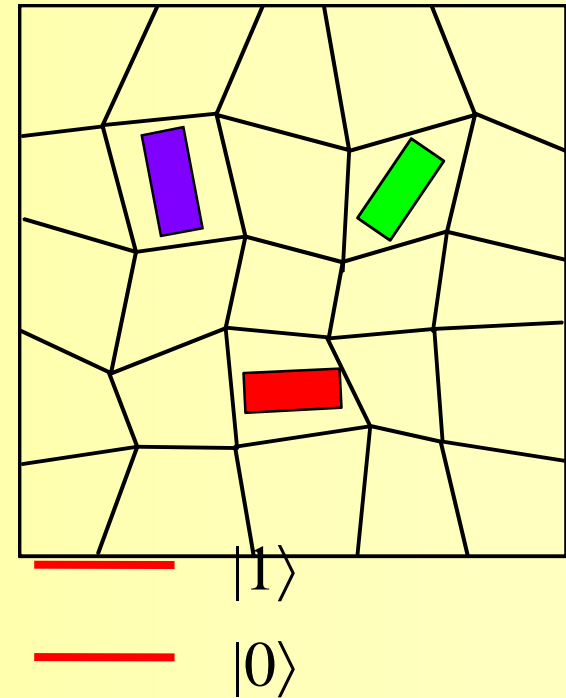
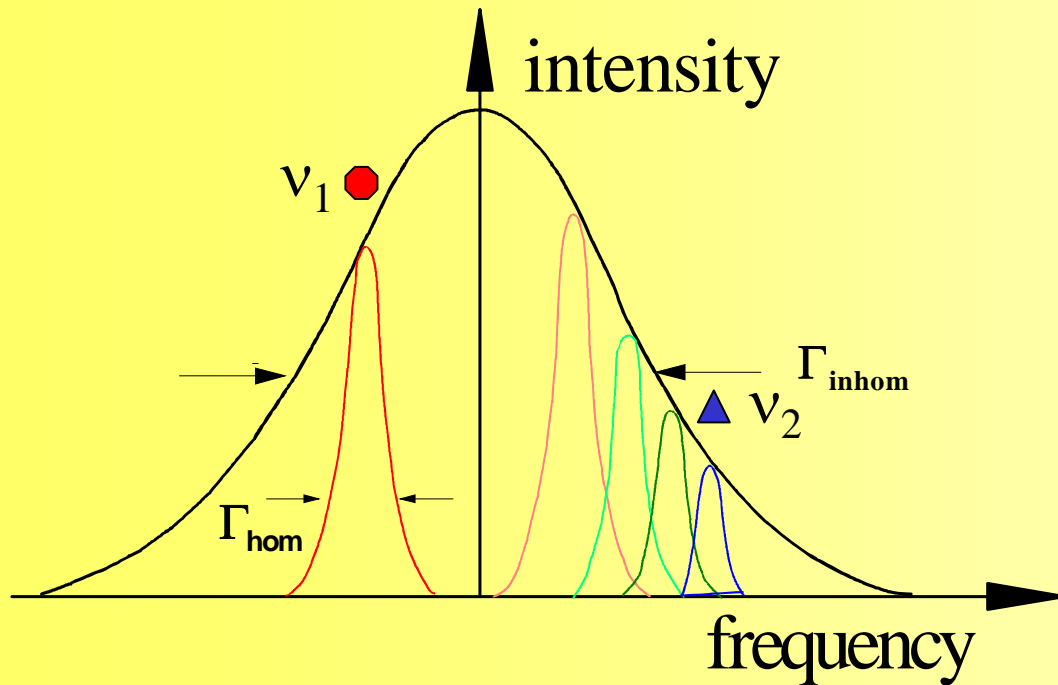


# Requirements for quantum computing

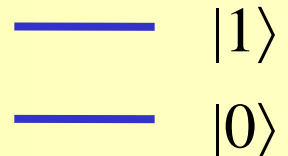
- Coherent two-level systems acting as qubits
- **Possibility to manipulate the qubits individually (single qubit operations)**
- Coupling between any two qubits (two-bit gates)
- Possibility for reliable read-out of the individual qubits
- Scalability

# 4f-4f absorption line from dopant ions in a rare earth doped inorganic crystal

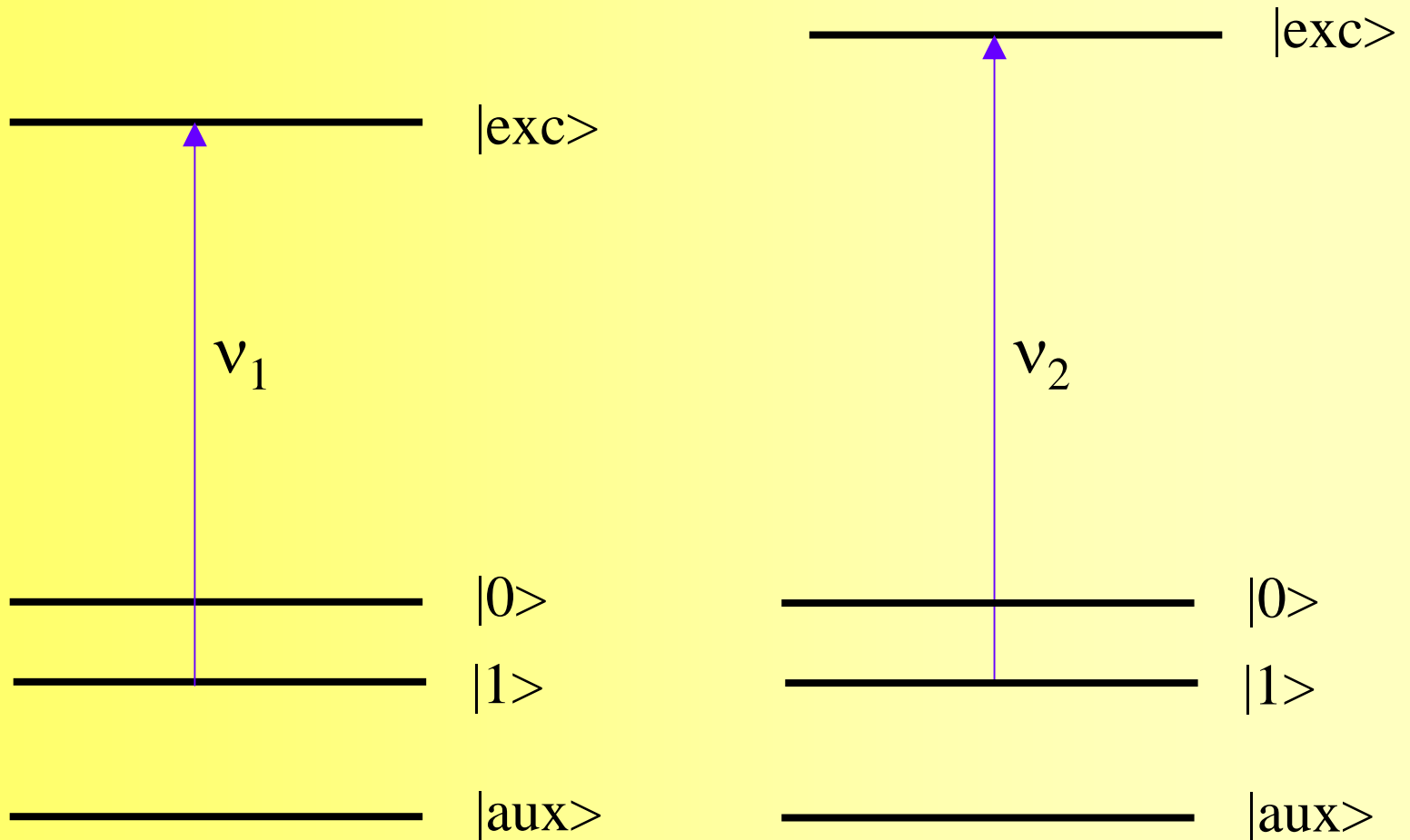
$\nu_1$  ●  $\nu_2$  ▲  
 Conceptual picture of crystal with dopant ions  $|exc\rangle$



- Narrow homogeneous line-widths (1-10 kHz)
- Large inhomogeneous line-widths (1-200 GHz)



# Addressing two different qubits in a rare-earth quantum computer



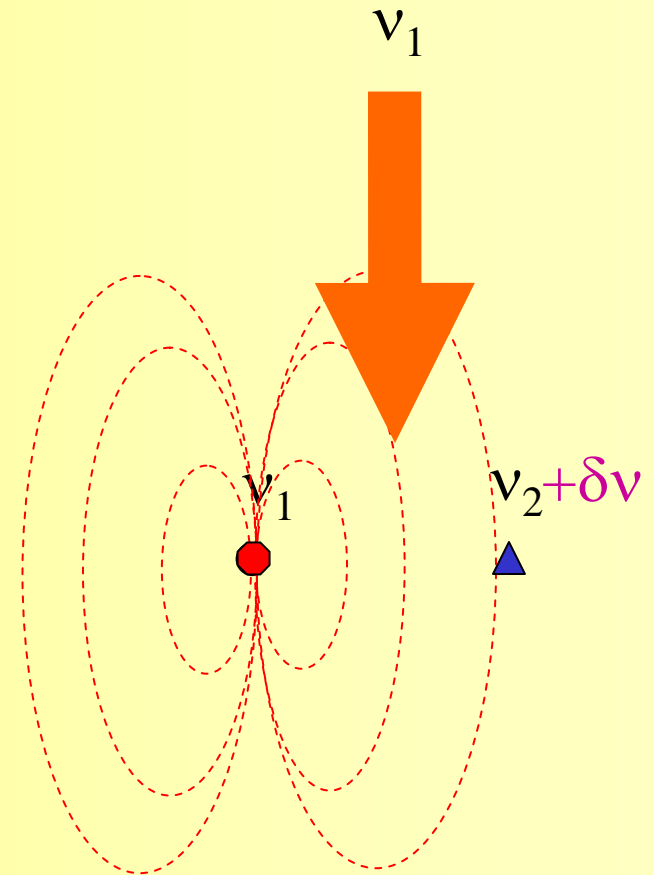
# Requirements for quantum computing

- Coherent two-level systems acting as qubits
- Possibility to manipulate the qubits individually (single qubit operations)
- **Coupling between any two qubits (two-bit gates)**
- Possibility for reliable read-out of the individual qubits
- Scalability



# Dipole-dipole interaction

1. Two ions absorbing at different frequencies are located close to each other in the crystal lattice. In a non-centrosymmetric site the ions will have a permanent electric dipole moment and ground and excited state dipole moments can be different
2. One of the ions is excited on its optical transition
3. This change in dipole moment is sensed by the other ion causing its absorption frequency to change.



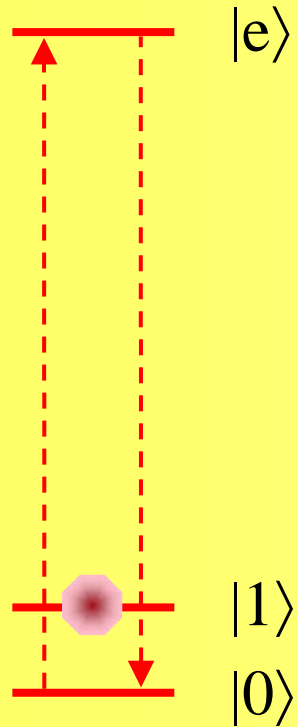
# Dipole-dipole interaction strength in rare-earth crystals

- Approximate numbers
- | <b>Ion distance</b> | <b>frequency shift</b> |
|---------------------|------------------------|
| • 100 nm            | 1 line width           |
| • 10 nm             | 1000 line widths       |
| • 1 nm              | 1000000 line widths    |

# Controlled-NOT quantum gate

Control

$v_1$  ●

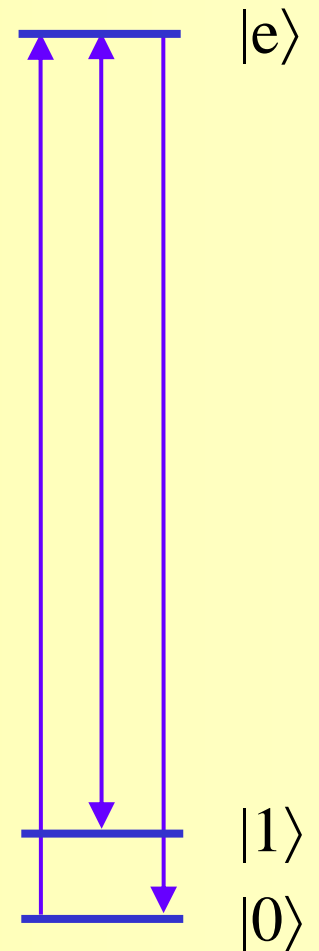


$v_1$  ●

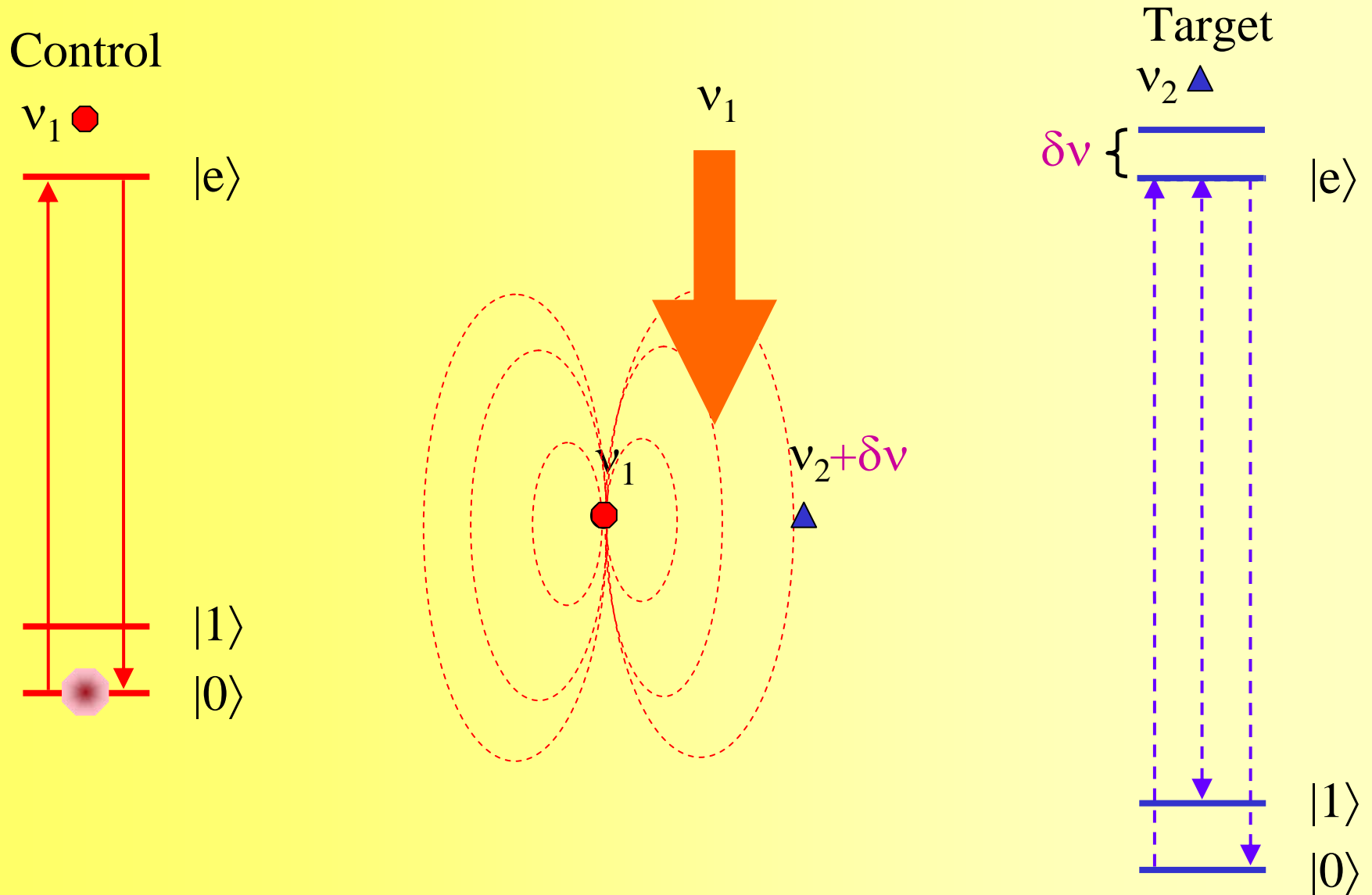
$v_2$  ▲

Target

$v_2$  ▲



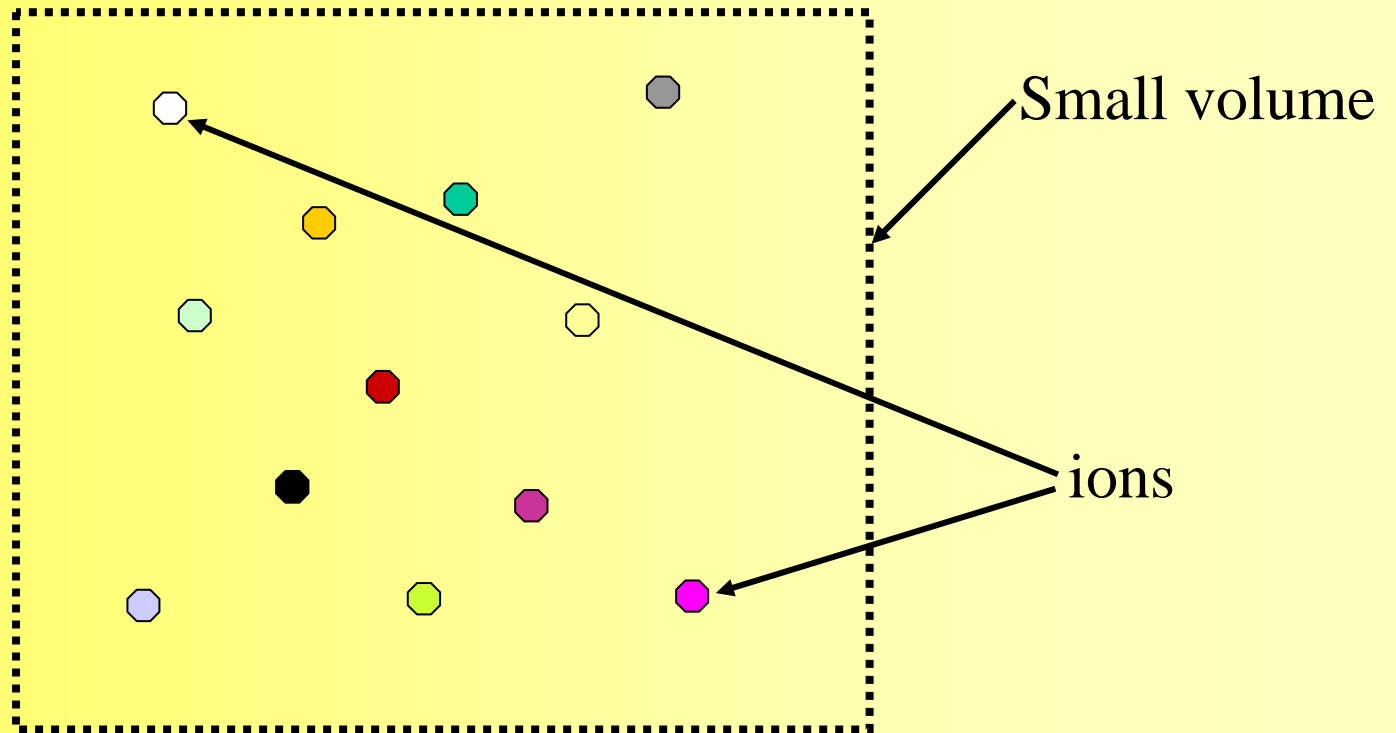
# Controlled-NOT quantum gate



# Requirements for quantum computing

- Coherent two-level systems acting as qubits
- Possibility to manipulate the qubits individually (single qubit operations)
- Coupling between any two qubits (two-bit gates)
- **Possibility for reliable read-out of the individual qubits**
- Scalability

All ions in a randomly selected small volume will interact strongly

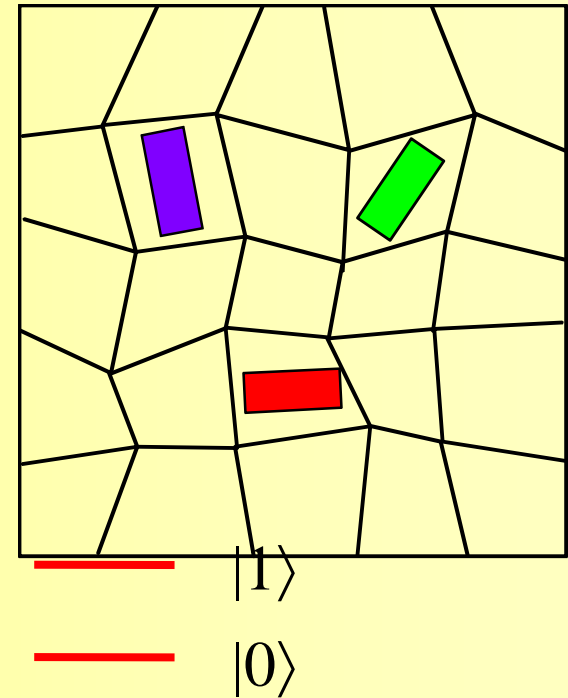
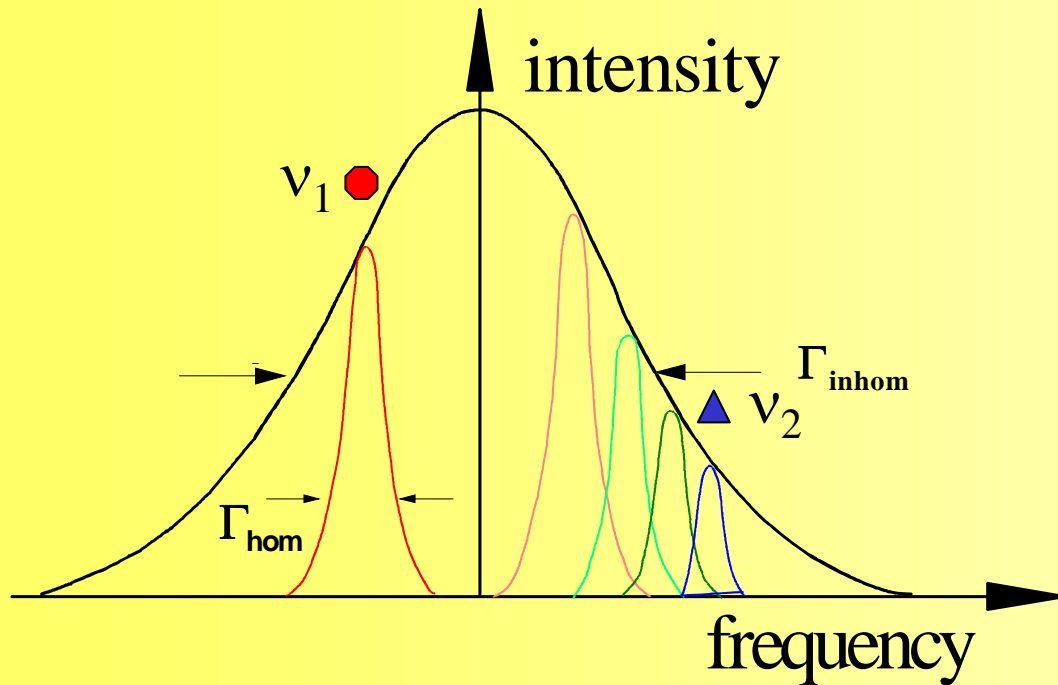


All ions interact strongly, but detecting single ions is difficult

# Absorption line from dopant ions in a rare earth doped inorganic crystal

Conceptual picture of crystal with dopant ions  $|exc\rangle$

$\nu_1$  ●  $\nu_2$  ▲



- Narrow homogeneous line-widths (1-10 kHz)
- Large inhomogeneous line-widths (1-200 GHz)

# Challenges with multiple instance approaches

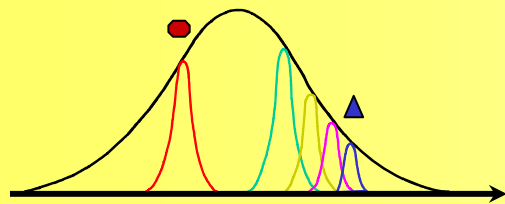
- All ions in a qubit (i.e. all QC instances) must have identical wave functions
- Both optical and hyperfine transitions are inhomogeneously broadened, ions in different instances may experience different laser field strengths and couple differently to the field
- Ion-ion interaction may differ between instances
- It may be possible to construct pulses which compensate for this but this adds to the operation time and increase decoherence



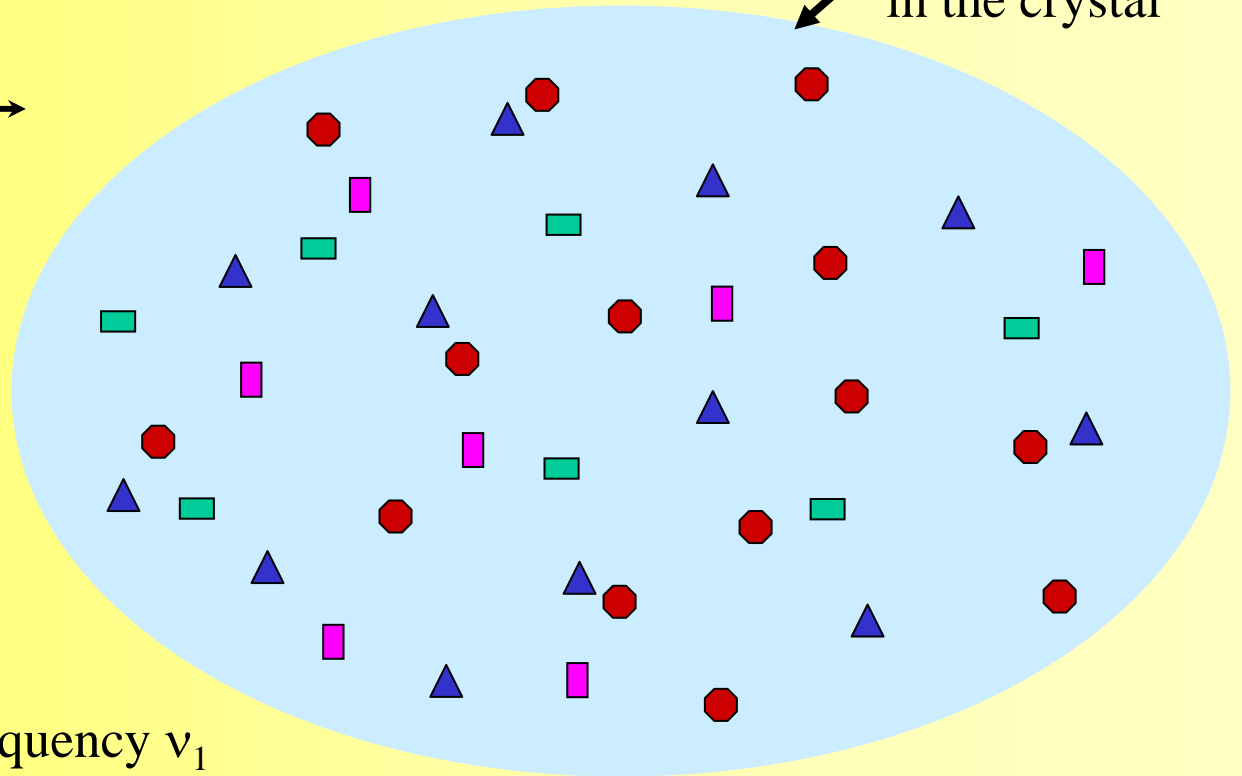
# Requirements for quantum computing

- Coherent two-level systems acting as qubits
- *Possibility to manipulate the qubits individually (single qubit operations)*
- *Coupling between any two qubits (two-bit gates)*
- **Possibility for reliable read-out of the individual qubits**
- Scalability

# Qubit distillation



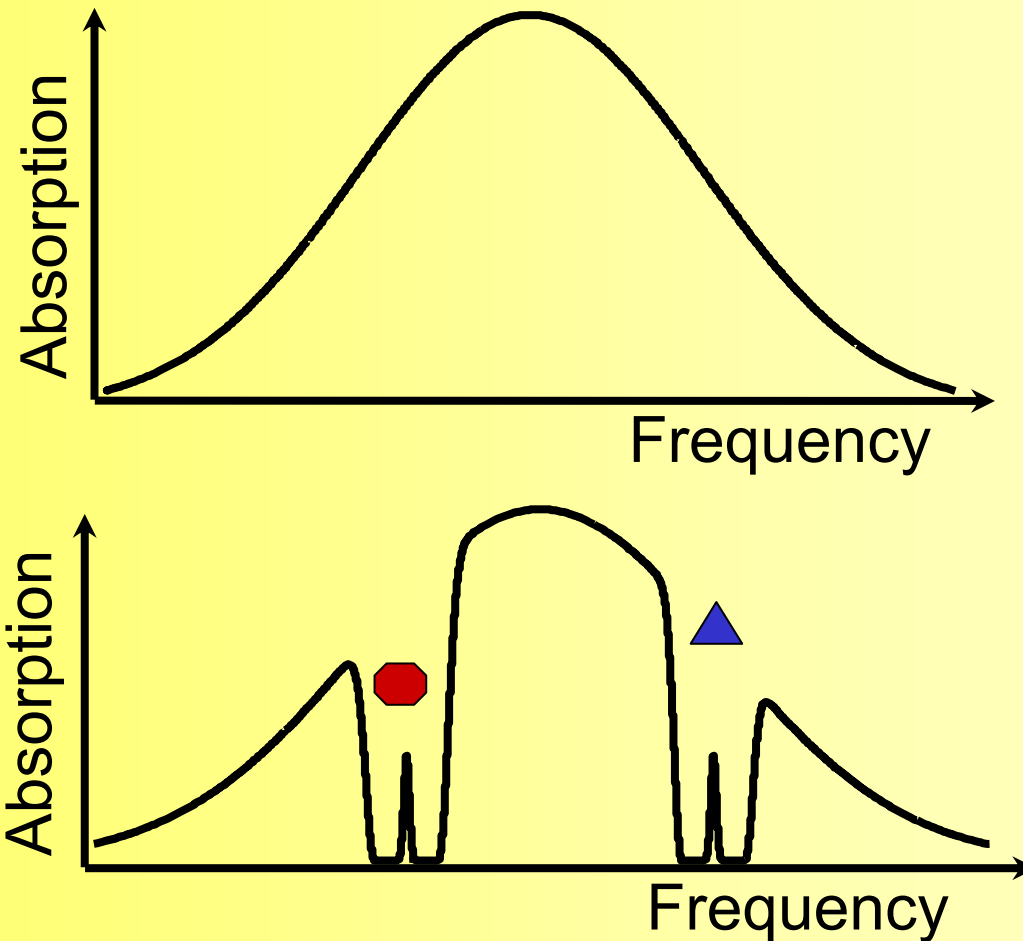
Arbitrary volume  
in the crystal



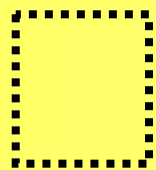
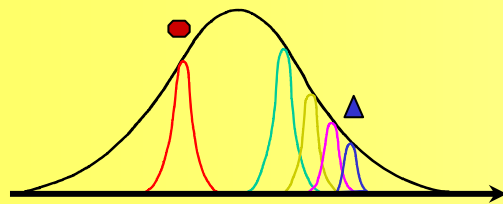
● ion absorbing at frequency  $\nu_1$

▲ ion absorbing at frequency  $\nu_2$

Inhomogeneous absorption profile is tailored to create qubit structures



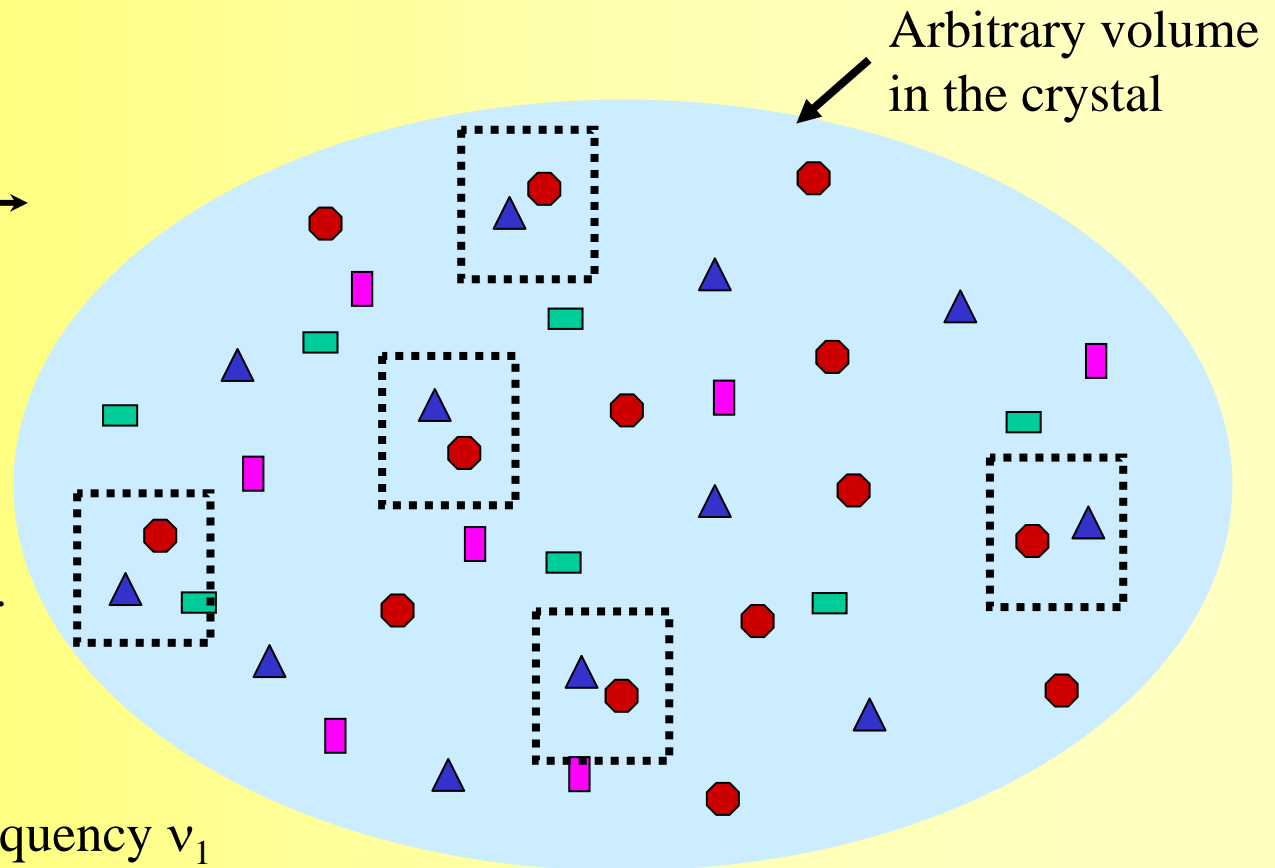
# Qubit distillation



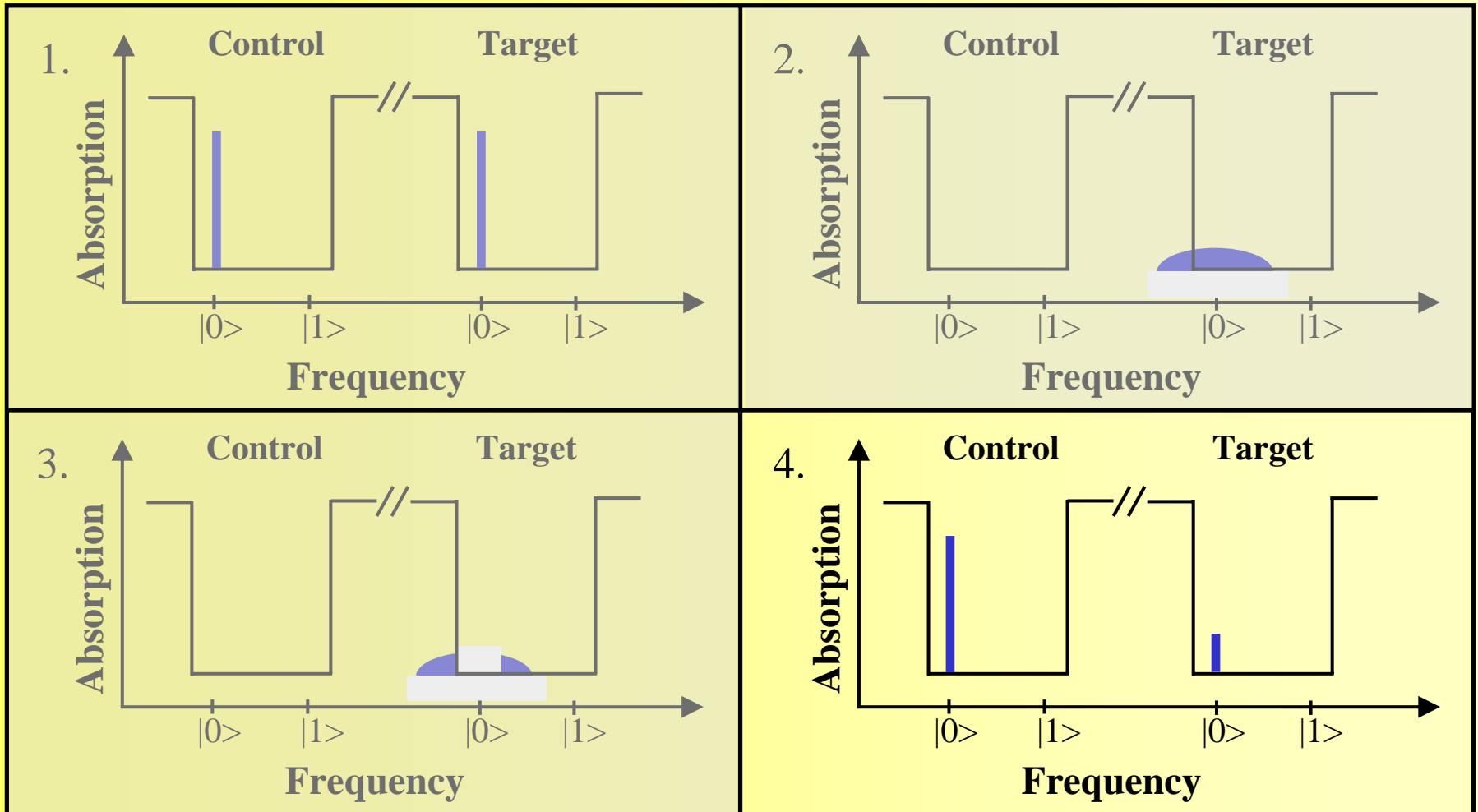
Ions interacting strongly enough for mutual control.  
Potential QC

● ion absorbing at frequency  $\nu_1$

▲ ion absorbing at frequency  $\nu_2$

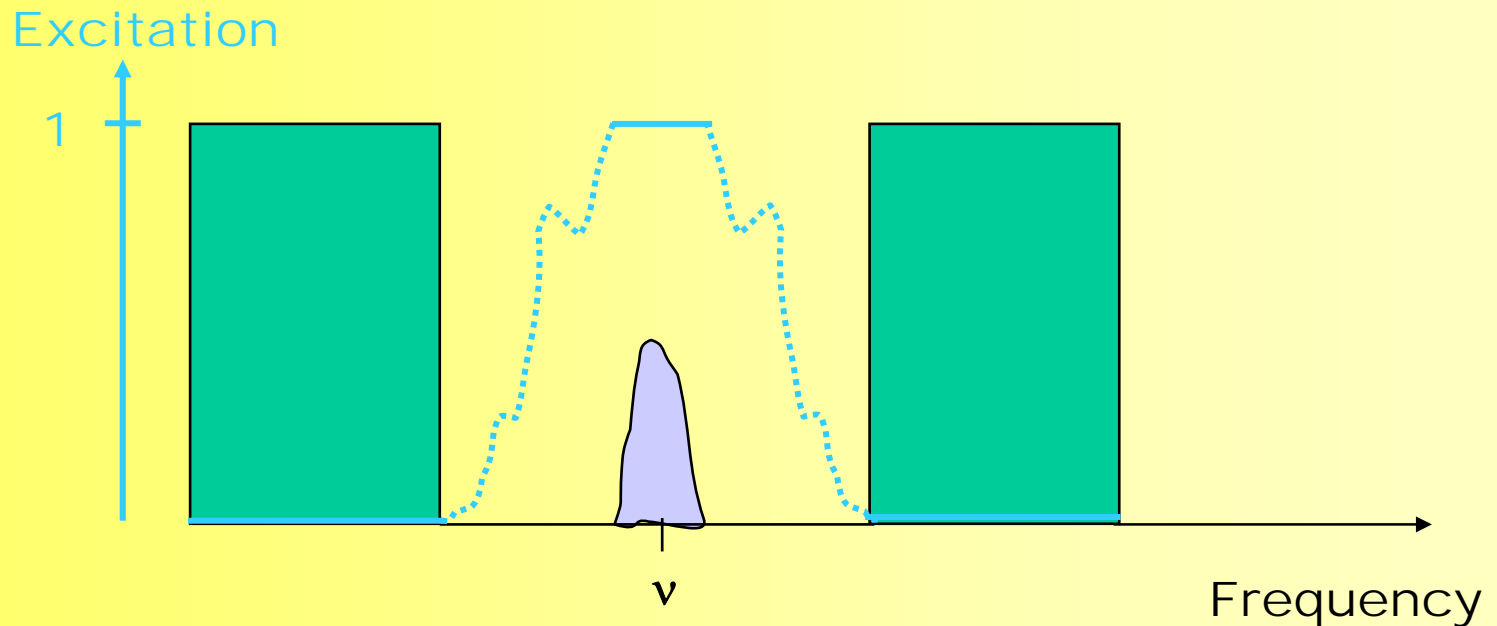


# Selecting strongly interacting ions

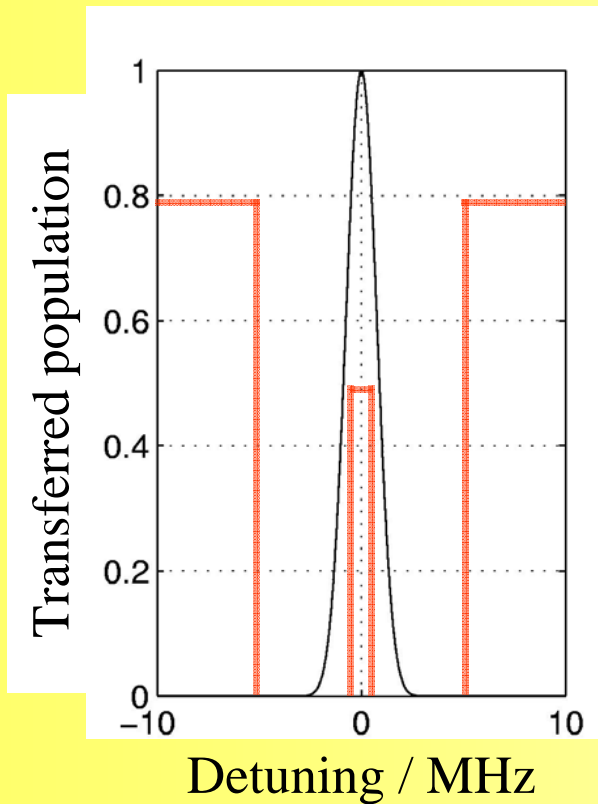


# Requirements on excitation pulses

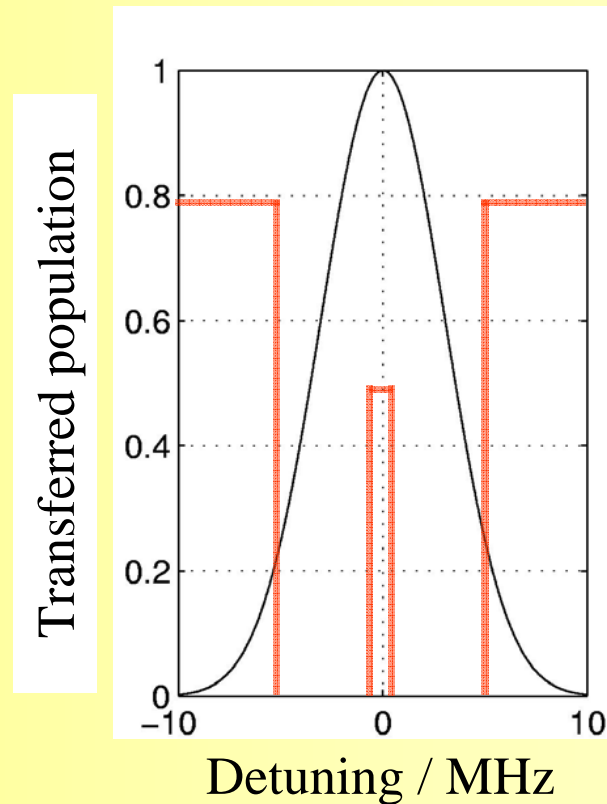
# How to interact with the qubit ions without interacting with ions at nearby absorption frequencies



# Excitation with Gaussian $\pi$ -pulses



$1 \mu\text{s}$  pulse



$0.25 \mu\text{s}$  pulse



# Pulse shapes for coherent transfer of population

This work was carried out by Ingela Roos together with Klaus Mølmer

- “Robust quantum computing with composite pulse and coherent population trapping”, Phys Rev A **69**, 22321 (2004)



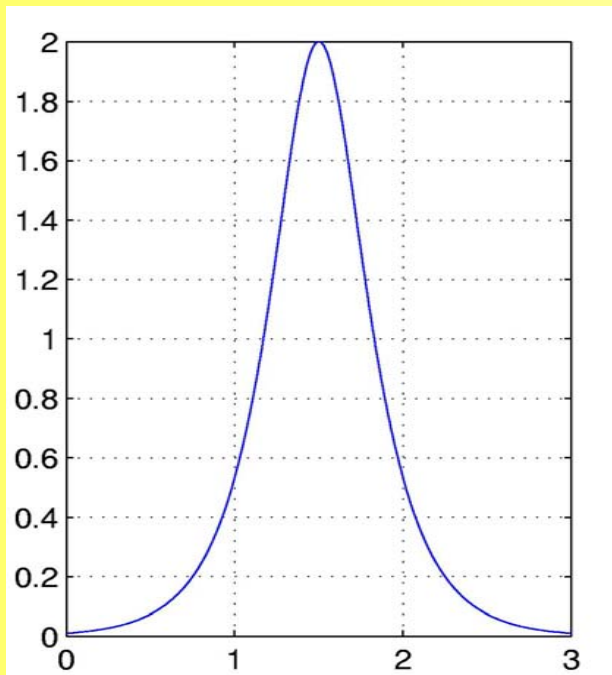
## Requirements

- Complete transfer of the peak of ions
- No excitation of surrounding ions

# Complex hyperbolic secant pulse

real sechyp  
amplitude envelope

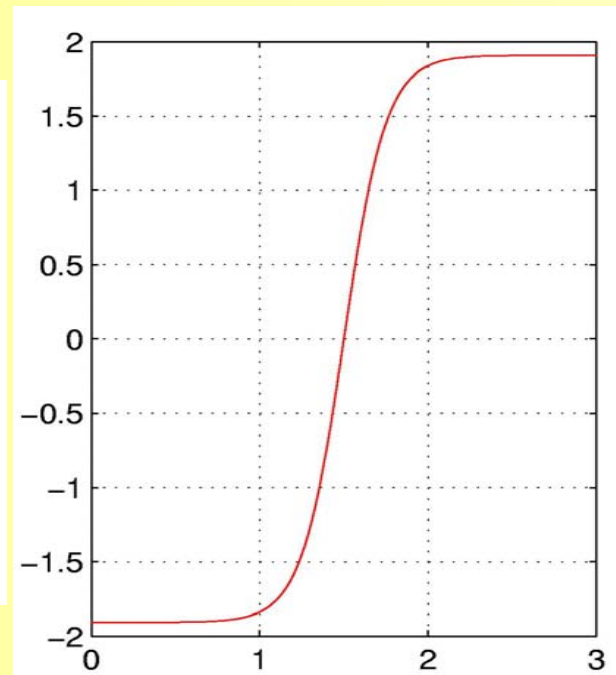
Rabi frequency / MHz



time /  $\mu\text{s}$

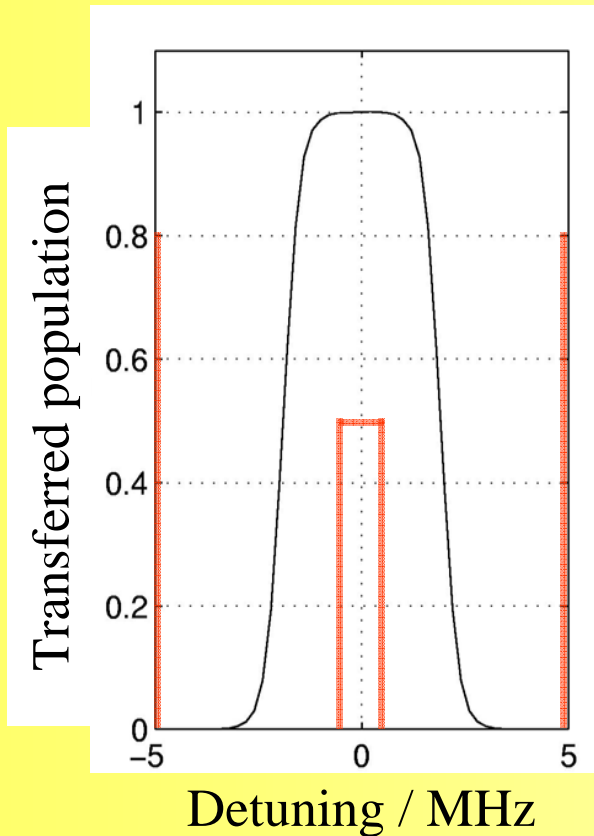
tanhyp frequency  
chirp

frequency / MHz

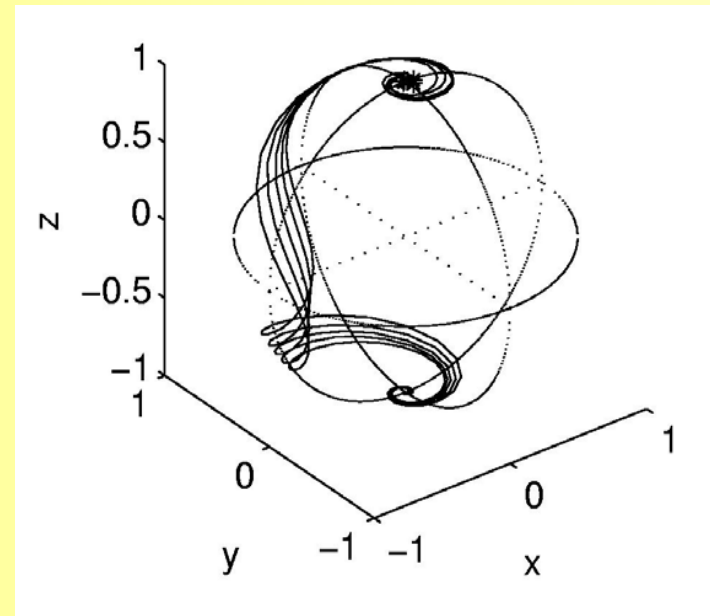


time /  $\mu\text{s}$

# Excitation with complex hyperbolic secant pulse



Evolution on the Bloch sphere

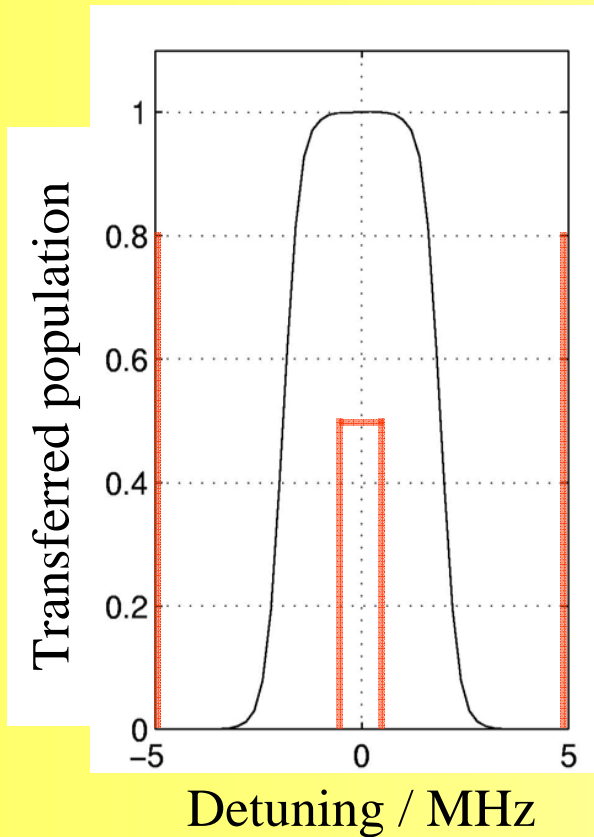


Further more, above a certain threshold intensity the operation is insensitive to different ions having different Rabi frequencies

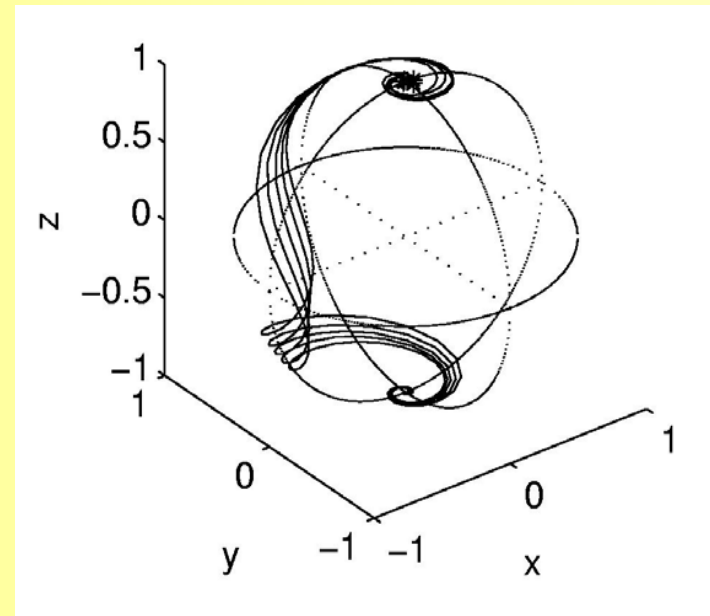
# We must have same wave function for all QC instances

- Hyperbolic secant pulses can compensate for
  - Finite channel widths
    - Variation in detunings
  - Field inhomogeneities & ion orientation variations
    - Variations in laser coupling

# Excitation with complex hyperbolic secant pulse



Evolution on the Bloch sphere



The right hand figure illustrates that the ions are driven coherently on the Bloch sphere

# High fidelity qubit operations require coherent laser systems

- Rare earth ion coherence times
  - Pr:YSO, optical  $\sim 100 \mu\text{s}$ , qubit  $\sim 100 \text{ ms}$
  - Eu:YSO, optical  $\sim 1 \text{ ms}$ , qubit ?
  - Er:YSO, optical ( $1.5 \mu\text{m}$ ) several ms
- Dye lasers
  - Commercial, coherence time  $< 1 \mu\text{s}$
  - Dortmund system coherence time  $\sim 10 \mu\text{s}$
  - Lund system coherence time  $\sim 100 \mu\text{s}$ 
    - Drift  $< 1 \text{ kHz/s}$

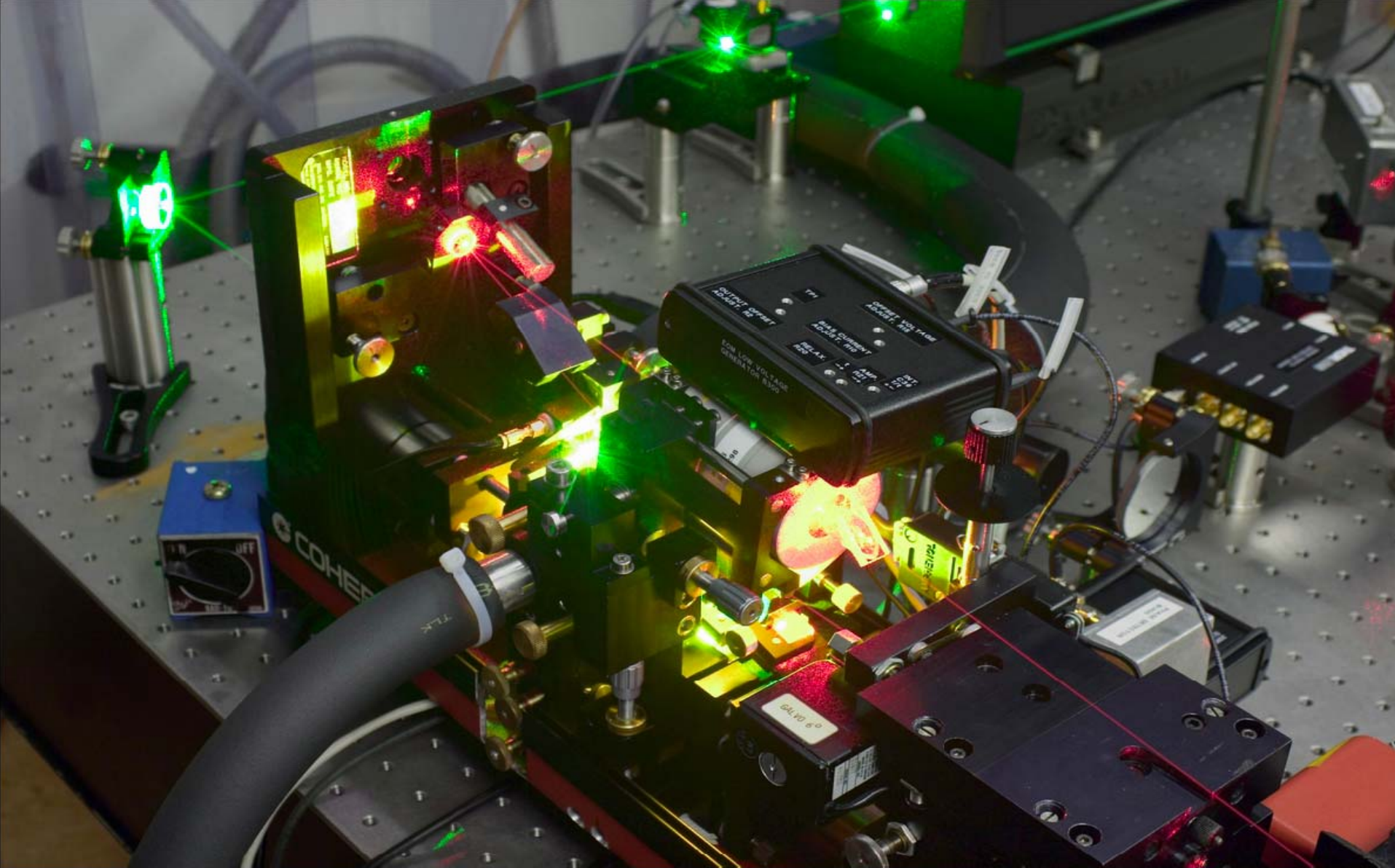


Photo: Tomas Svensson

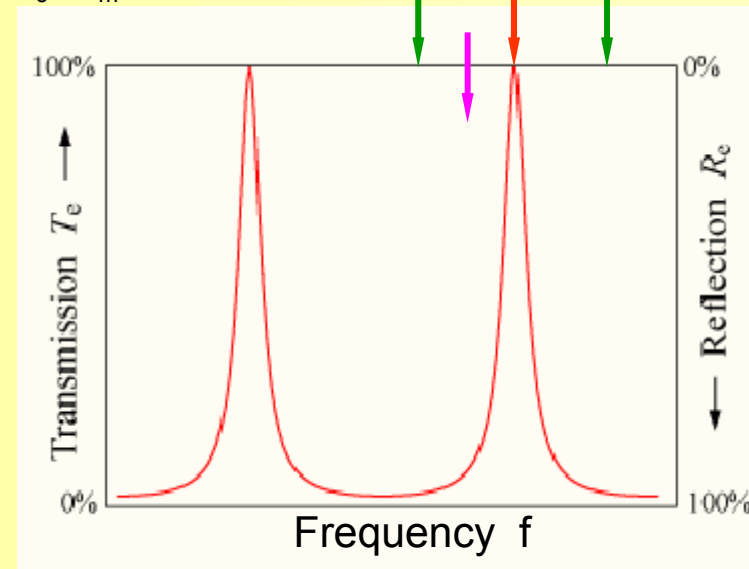
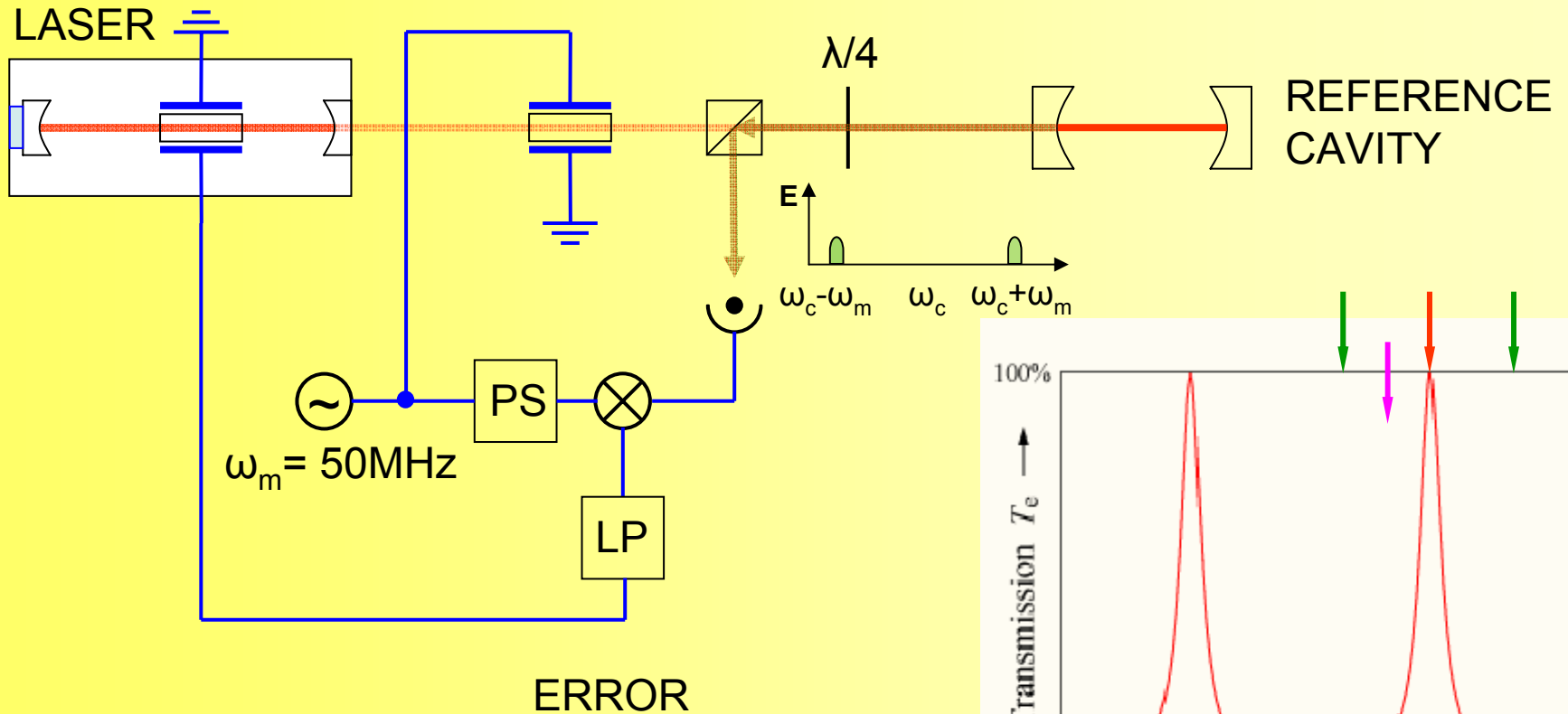
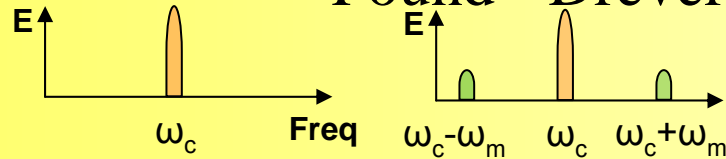


**Photo:**  
**Tomas Svensson**



# Phase stabilization against a cavity

## Pound - Drever - Hall

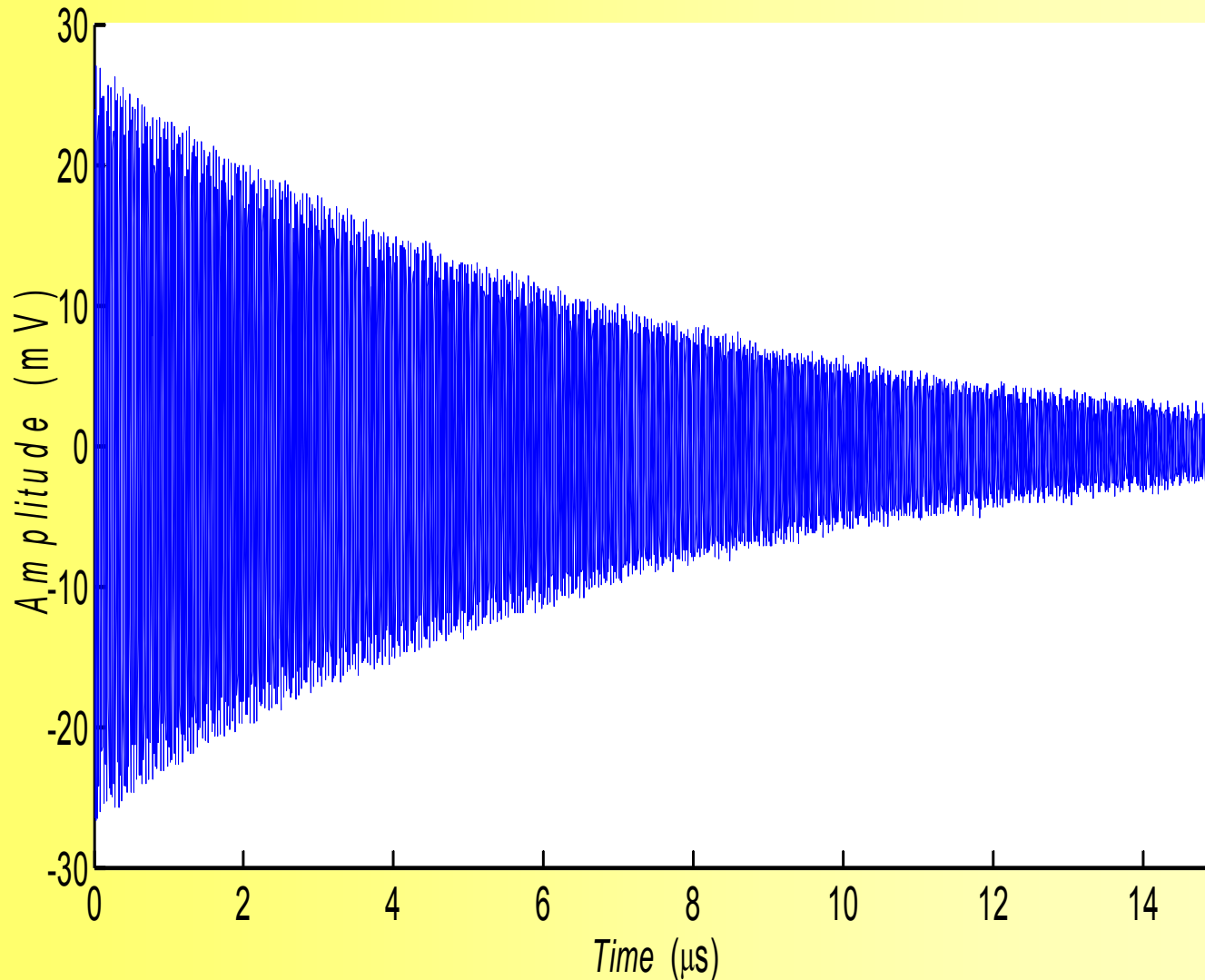


# Locking to spectral hole

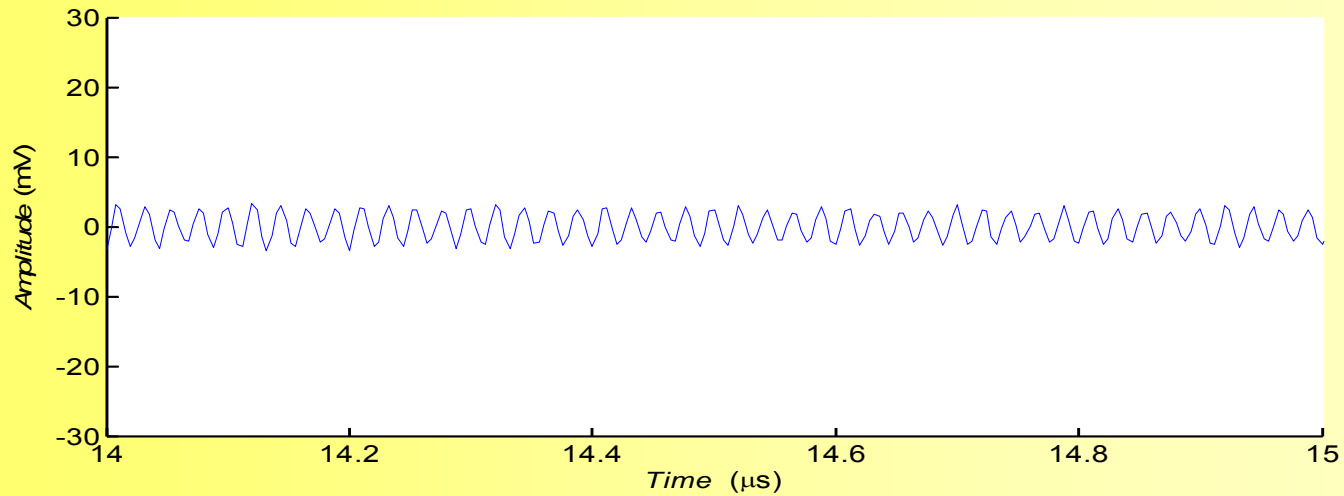
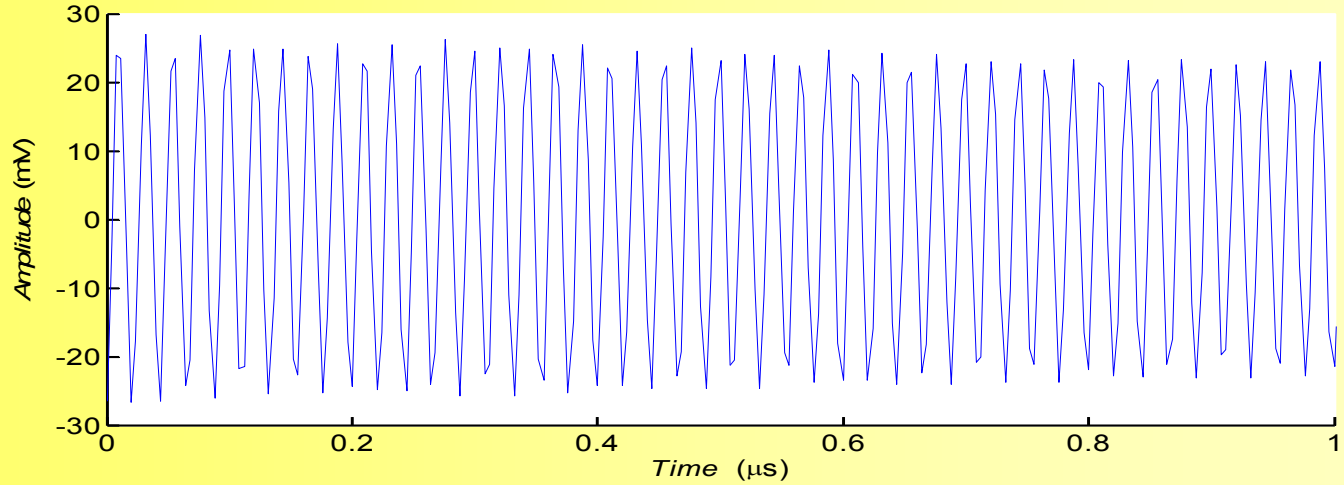
- Cavity with fixed transmission linewidth is replaced against a spectral hole with a linewidth that dynamically adapts to the laser linewidth
- Different optimum modulation index
- Optimum absorption
- New locking regimes
  - system locks leading to constant drift

# Free induction decay Pr:YSO

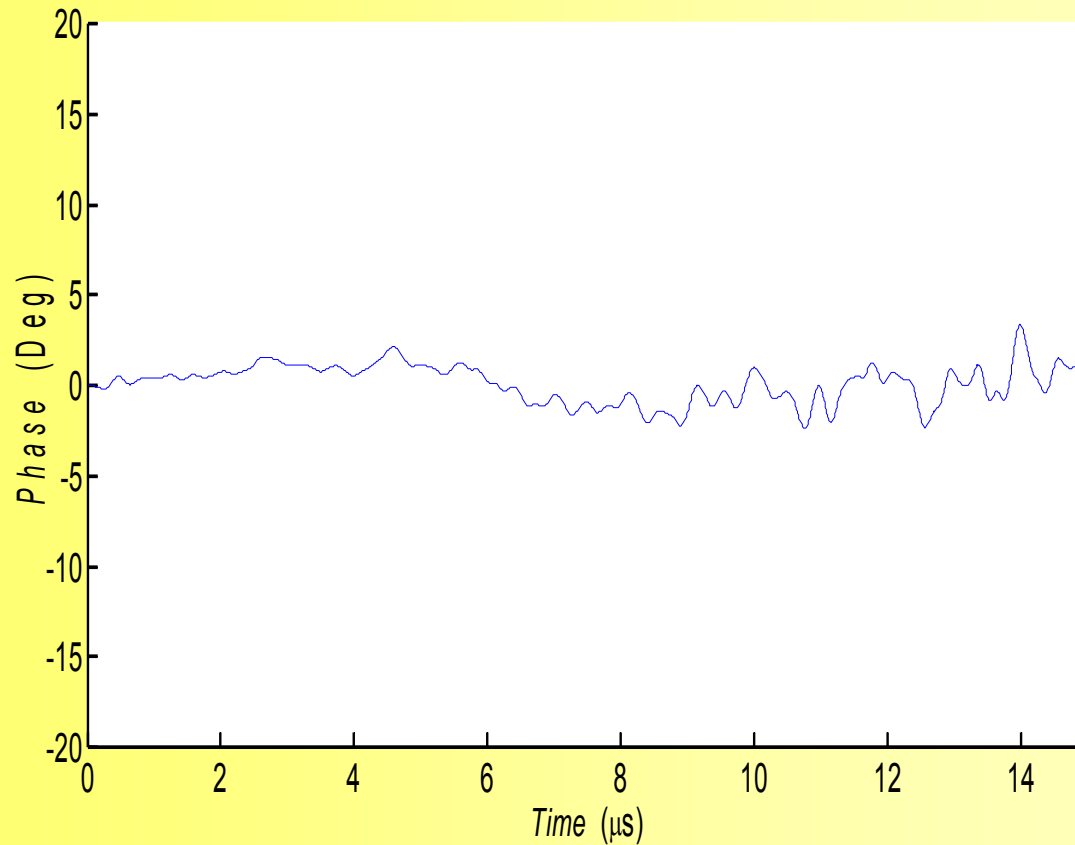
## Beat signal between laser and Pr ions



# Free induction decay

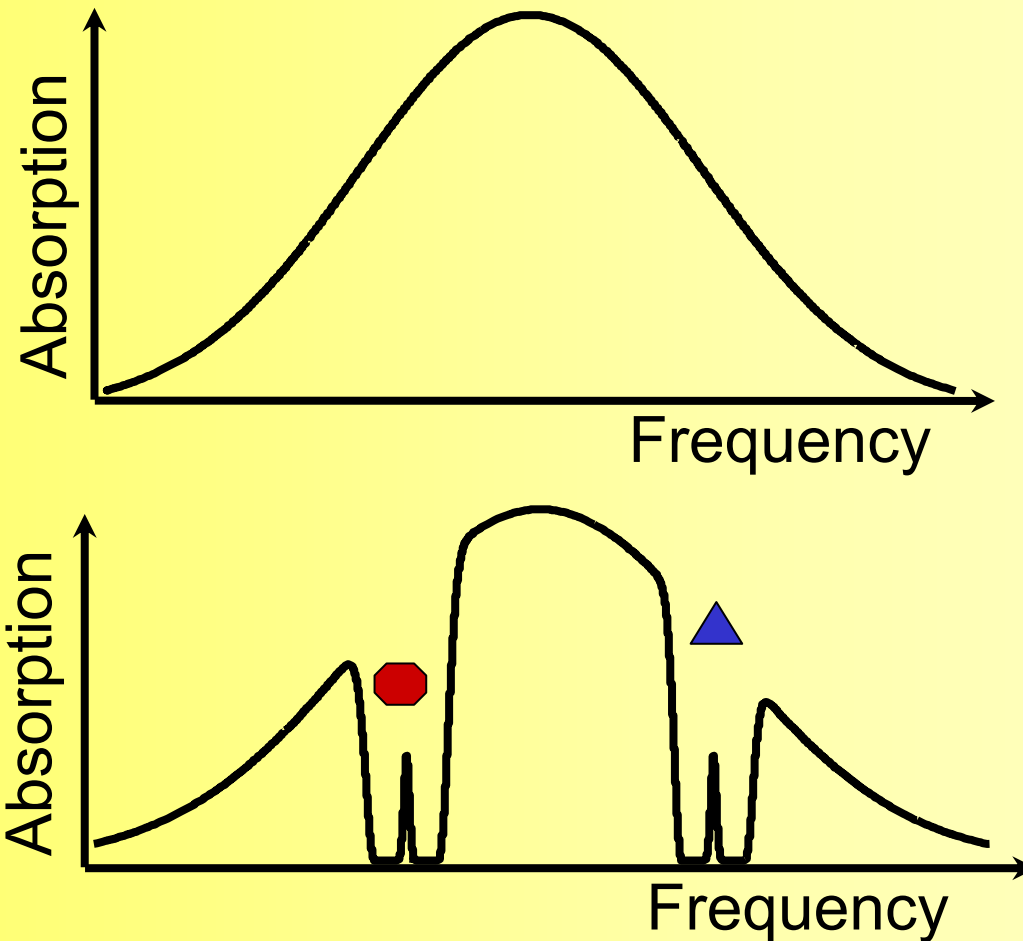


Laser phase drift  $< 5^\circ$  over  $10 \mu\text{s}$   
Coherence time  $> 100 \mu\text{s}$

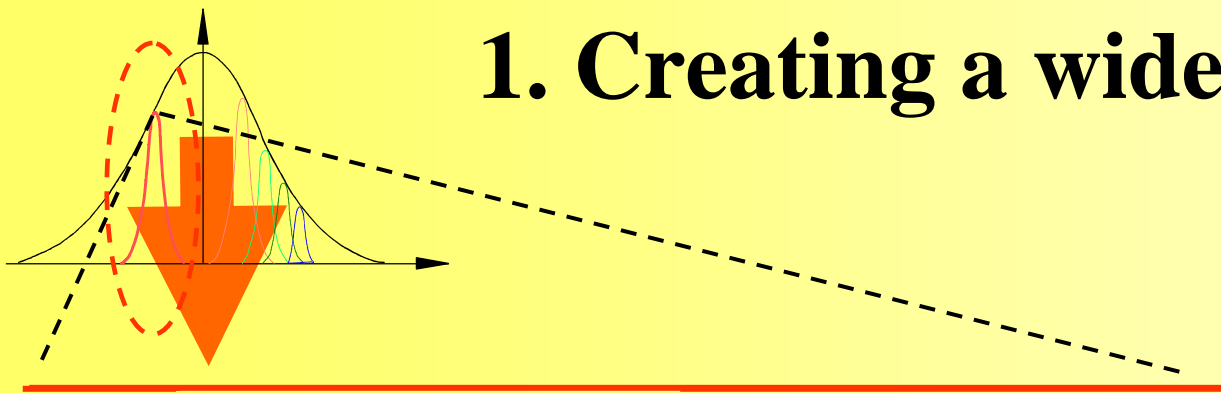


# Qubit creation

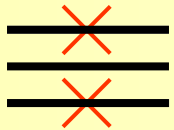
Inhomogeneous absorption profile is tailored to create qubit structures

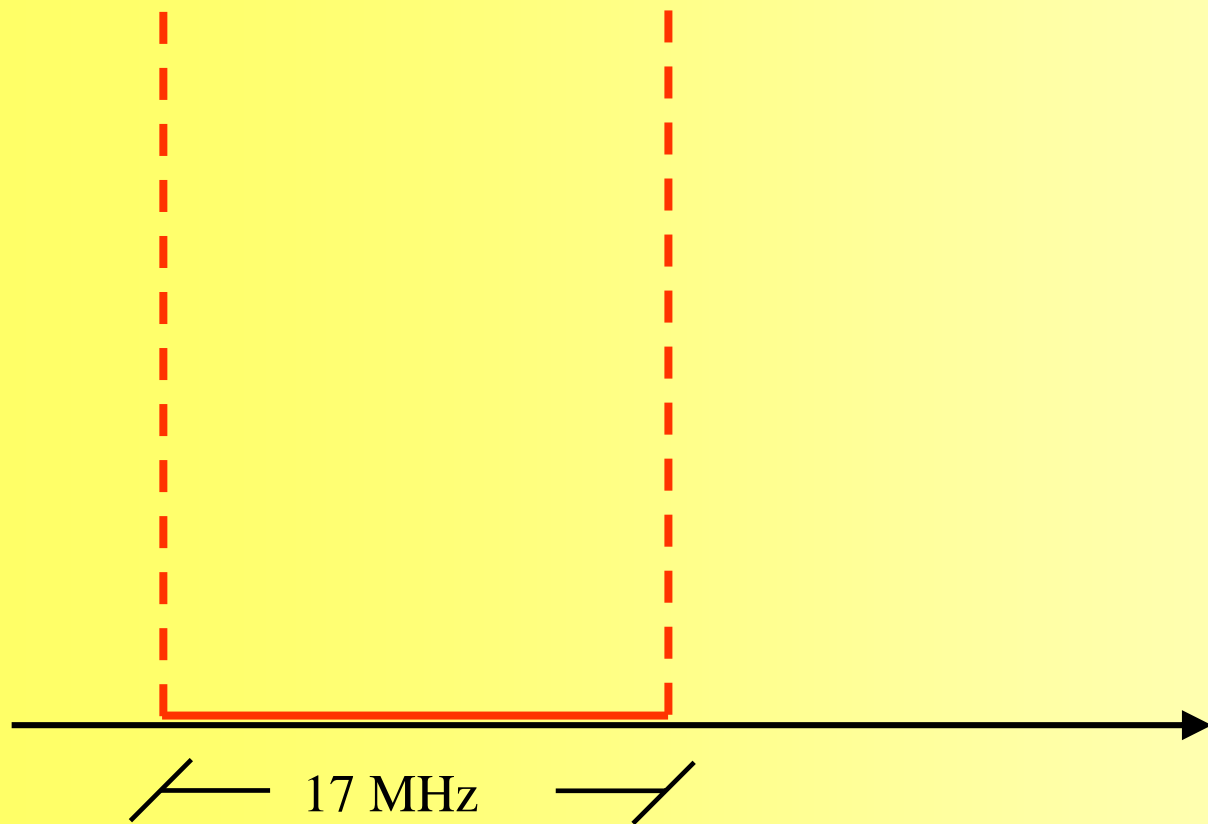


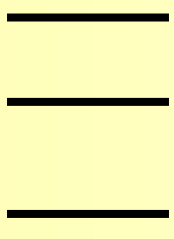
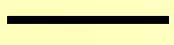
# 1. Creating a wide pit



Pr:Y<sub>2</sub>SiO<sub>5</sub>

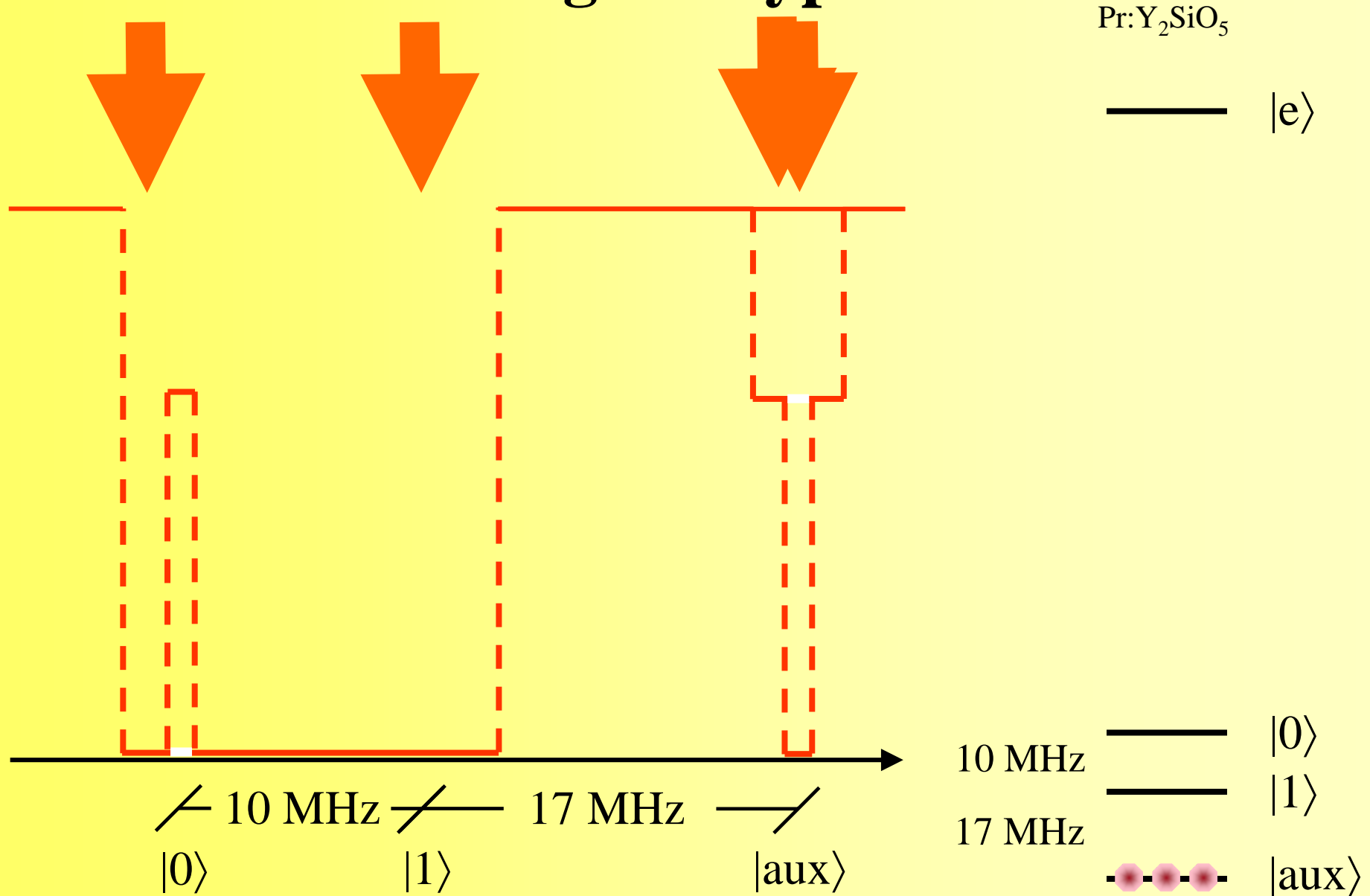
4.8 MHz  |e⟩  
4.6 MHz



10.2 MHz  |0⟩  
|1⟩  
17.3 MHz  |aux⟩

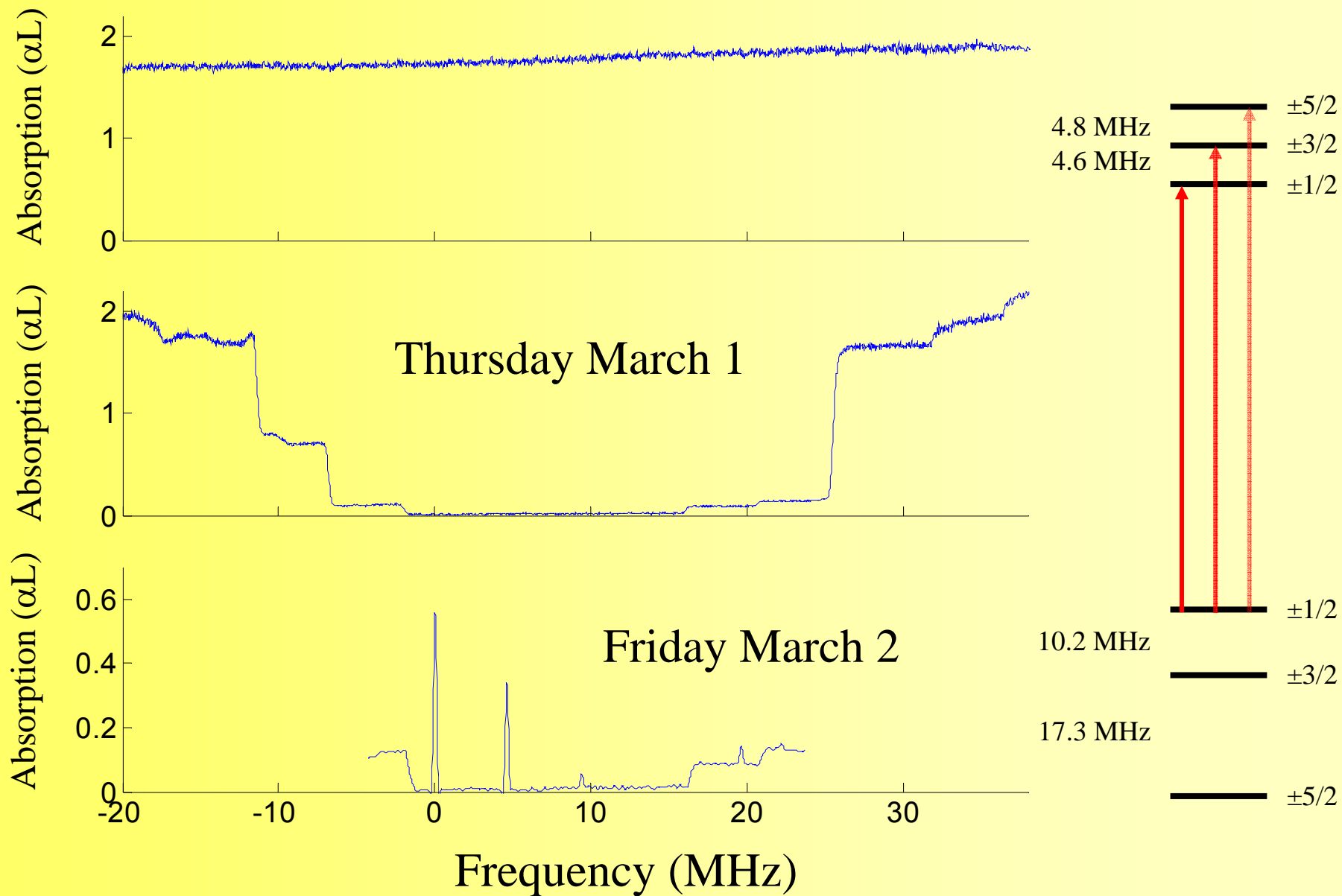


## 2. Isolating one type of ions

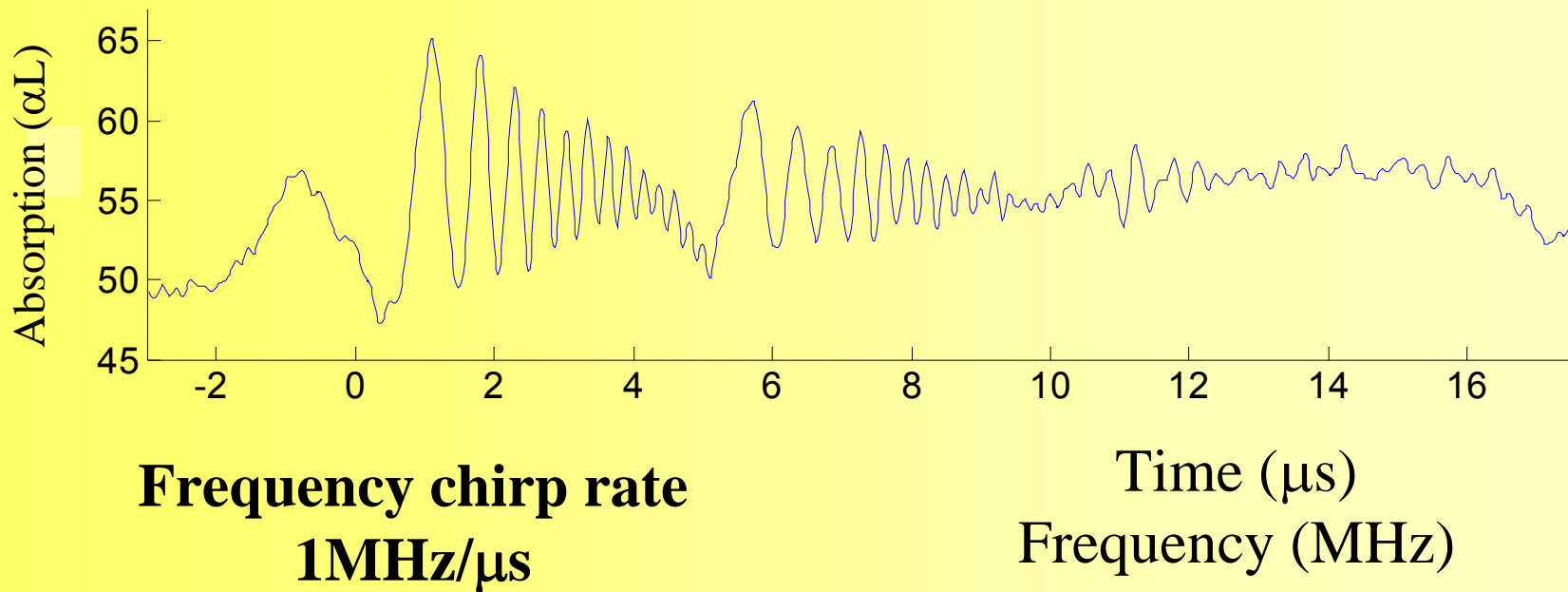
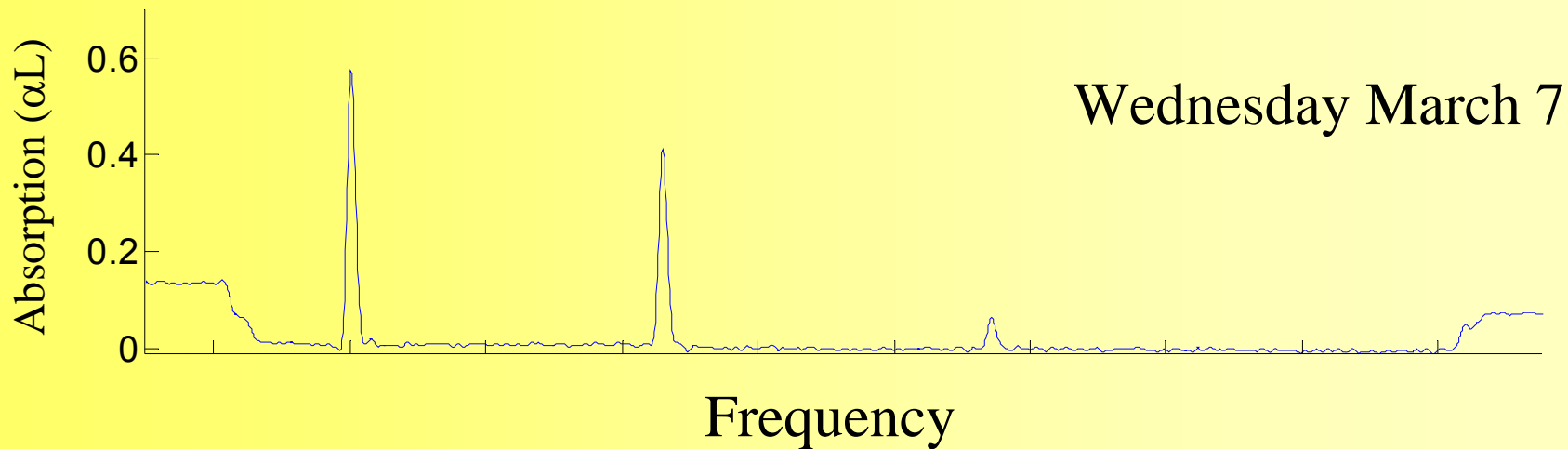


Experimental data

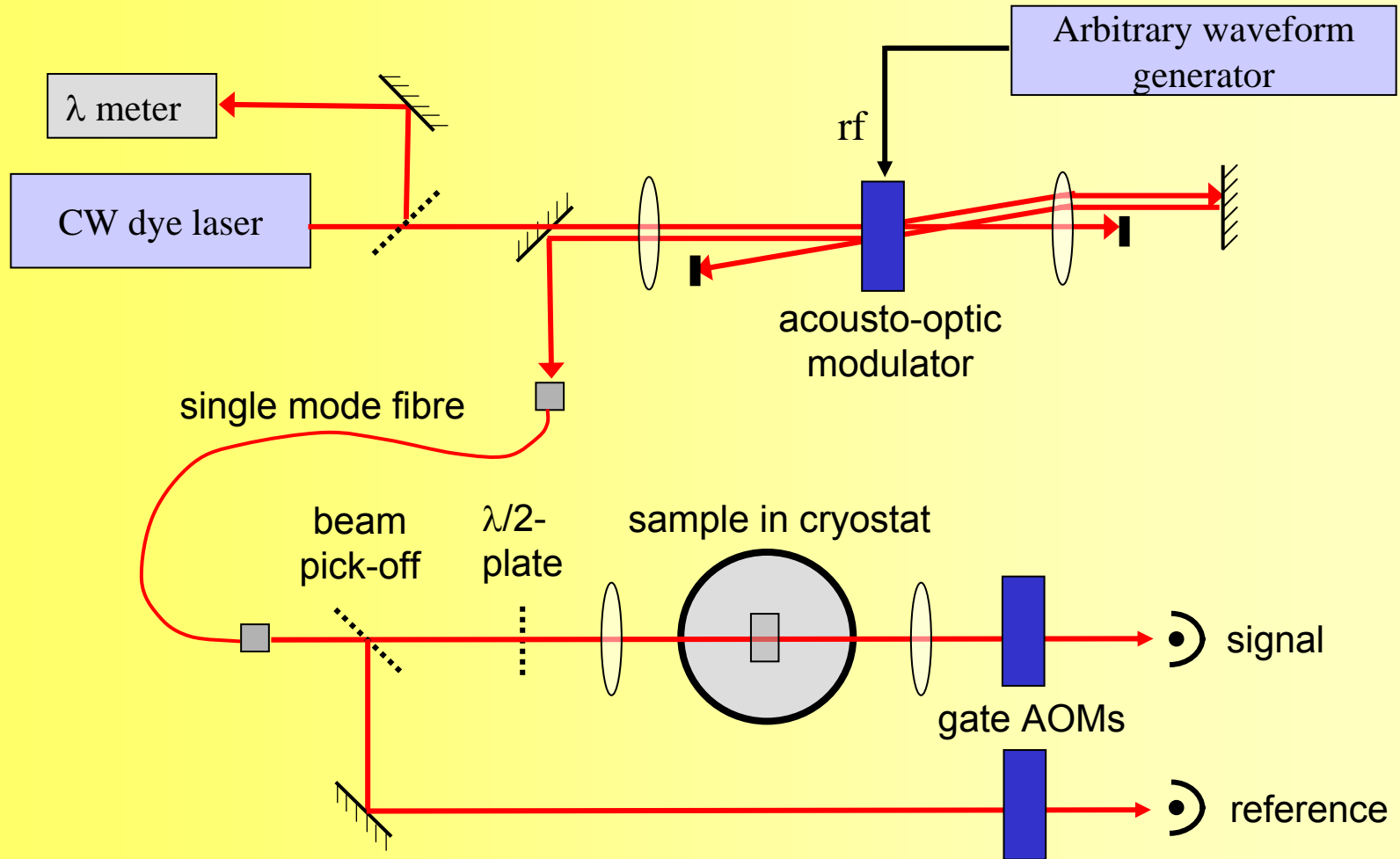
# Pit & peak creation



# Readout procedure



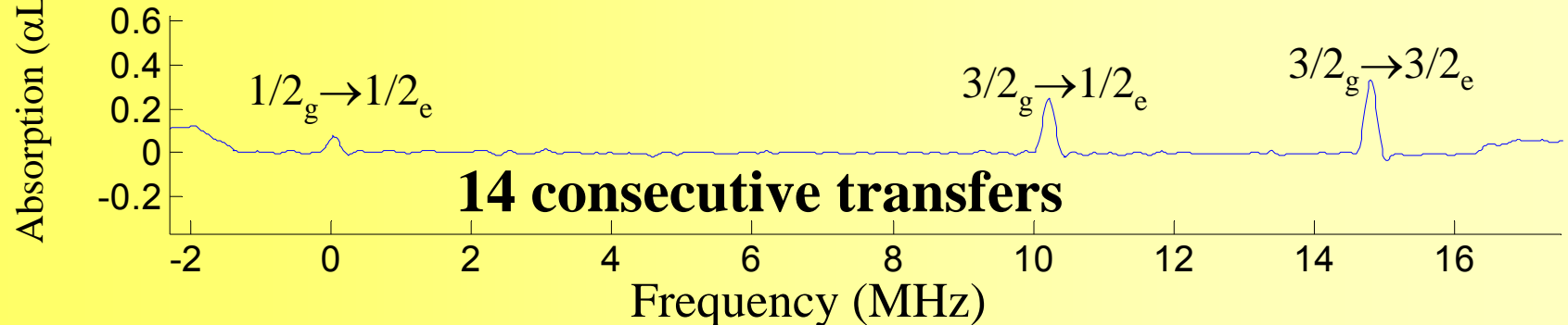
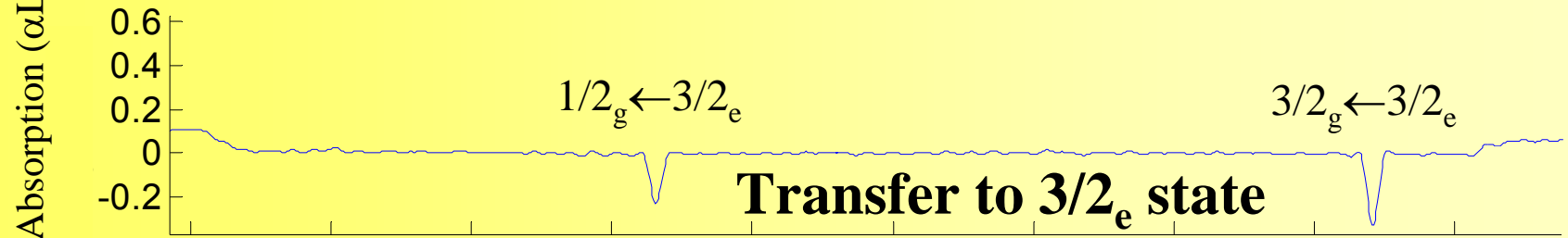
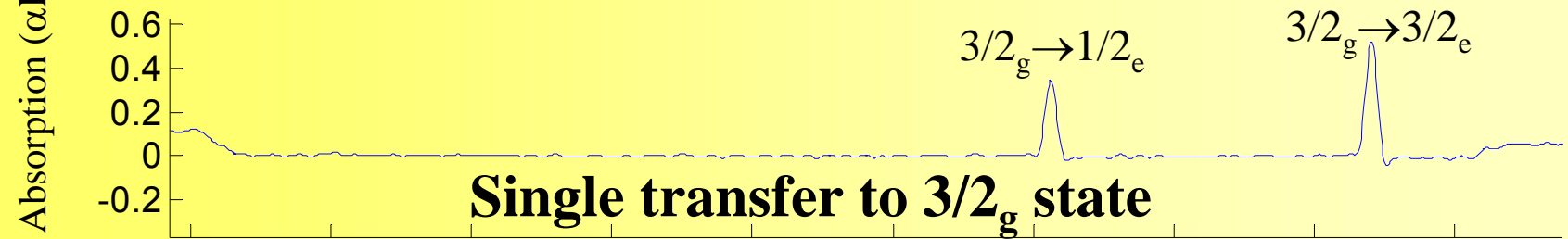
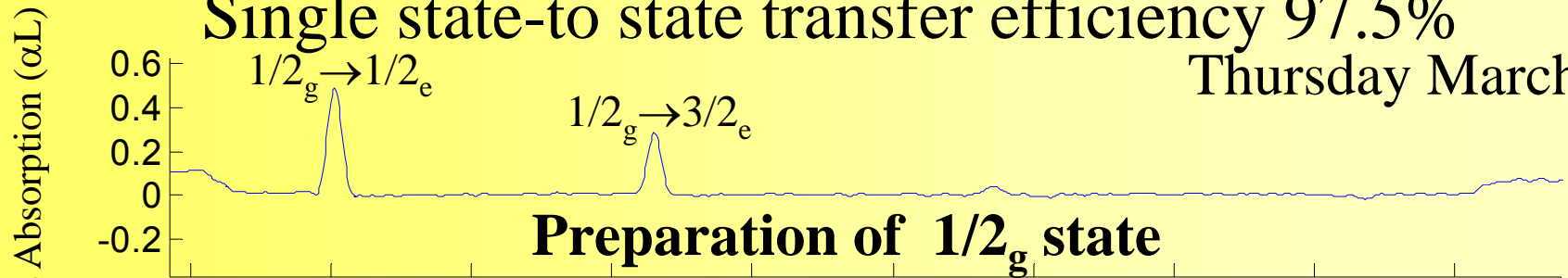
# Experimental set-up



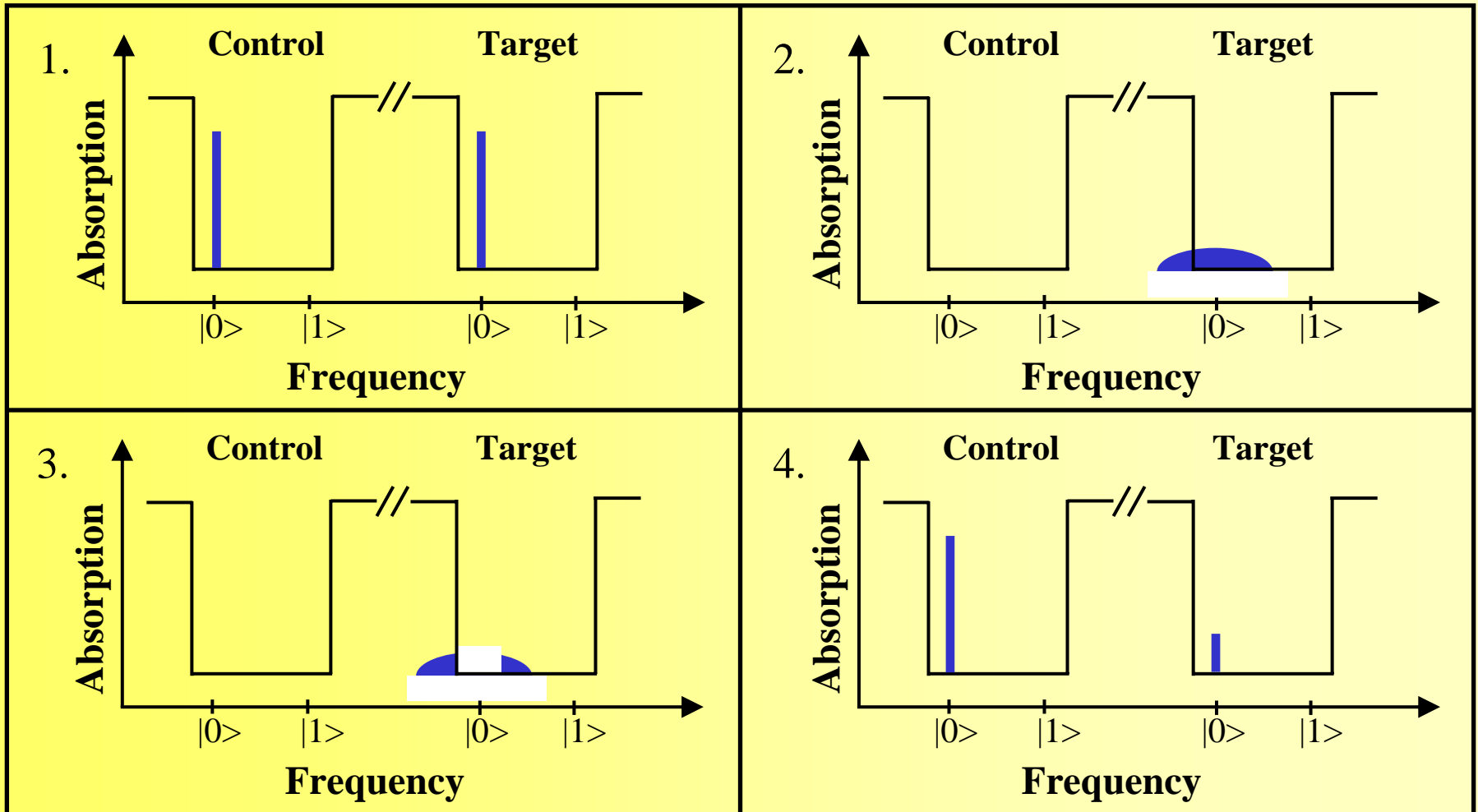
# State to state transfer

Single state-to state transfer efficiency 97.5%

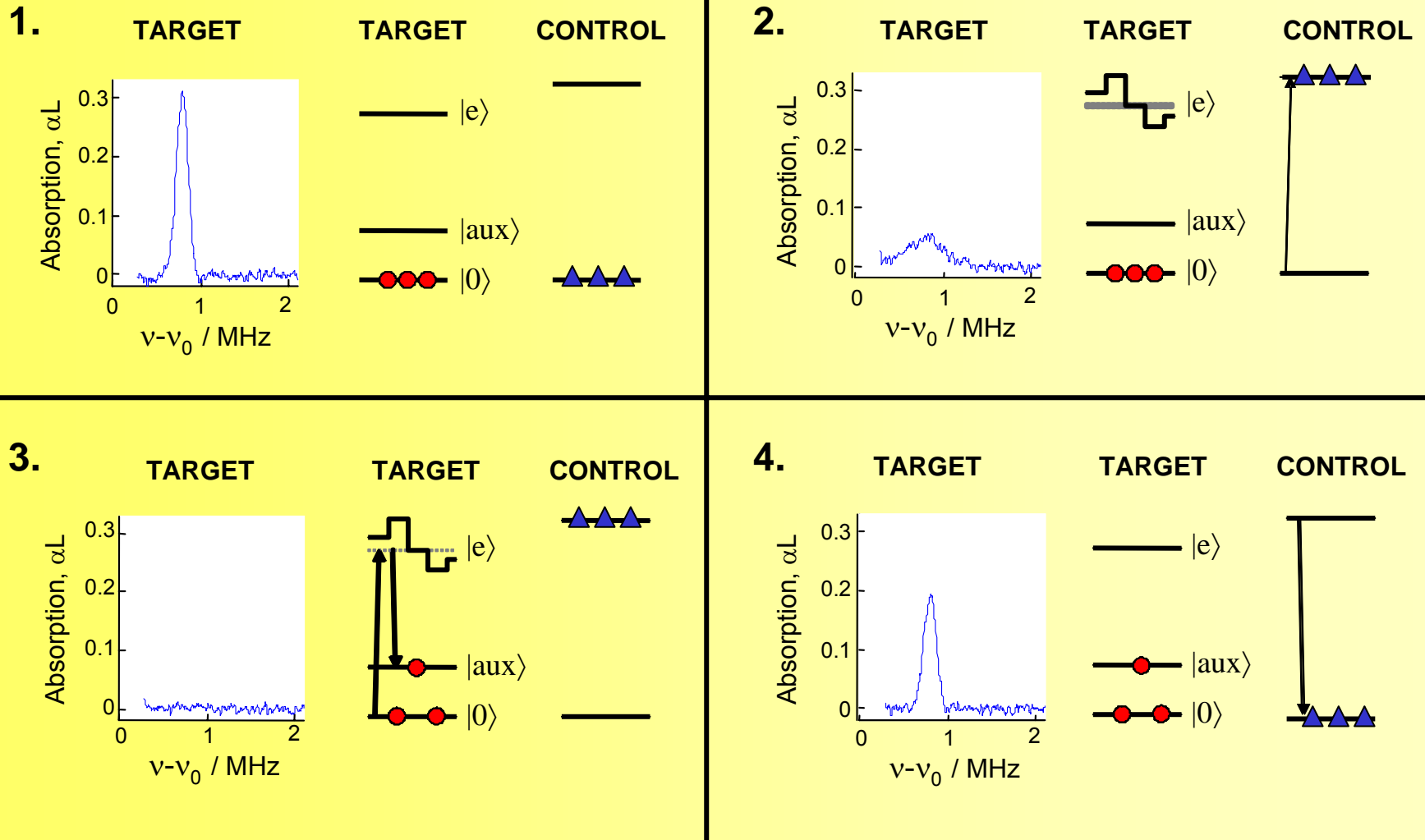
Thursday March 8



# Selecting strongly interacting ions

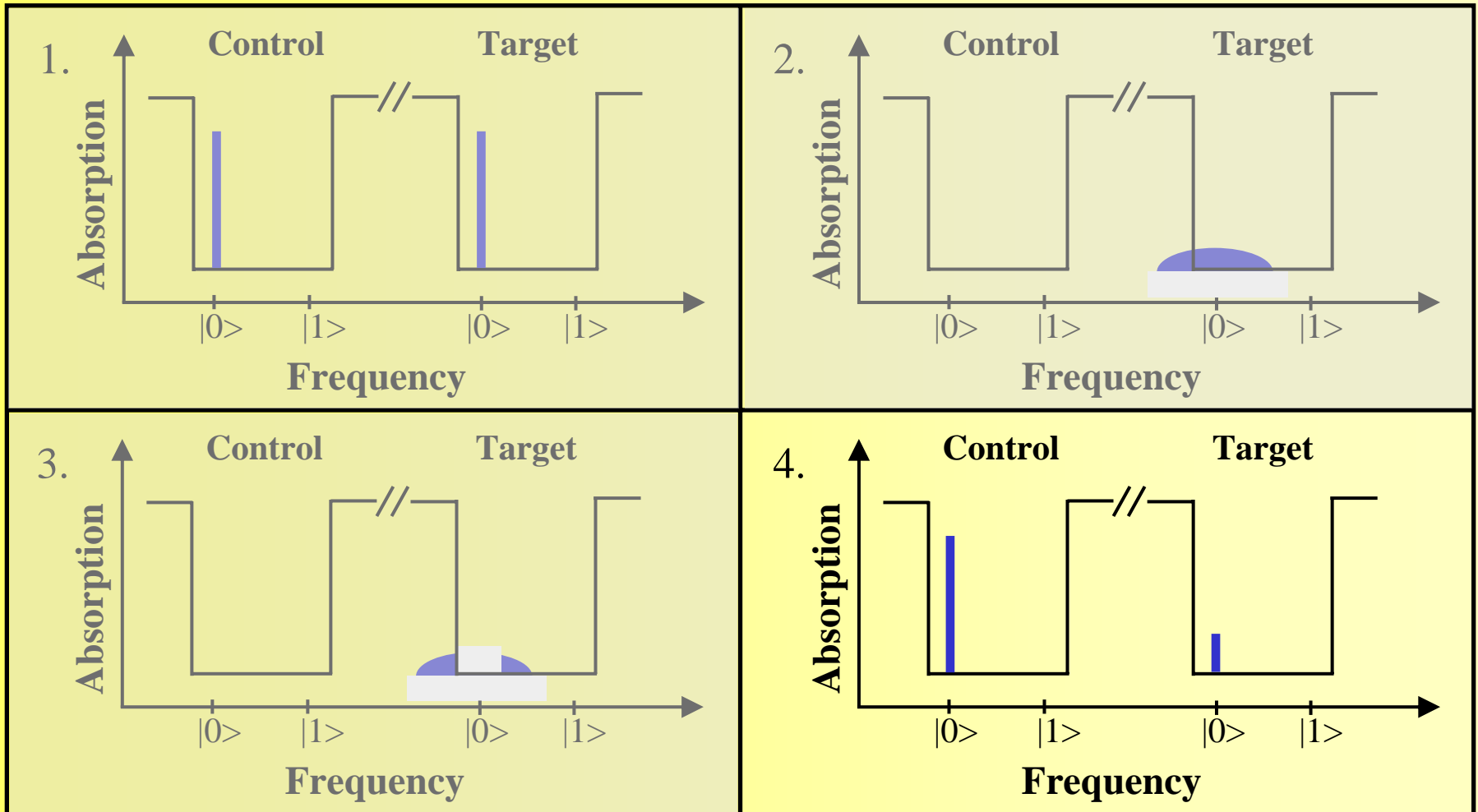


# Qubit distillation experiment





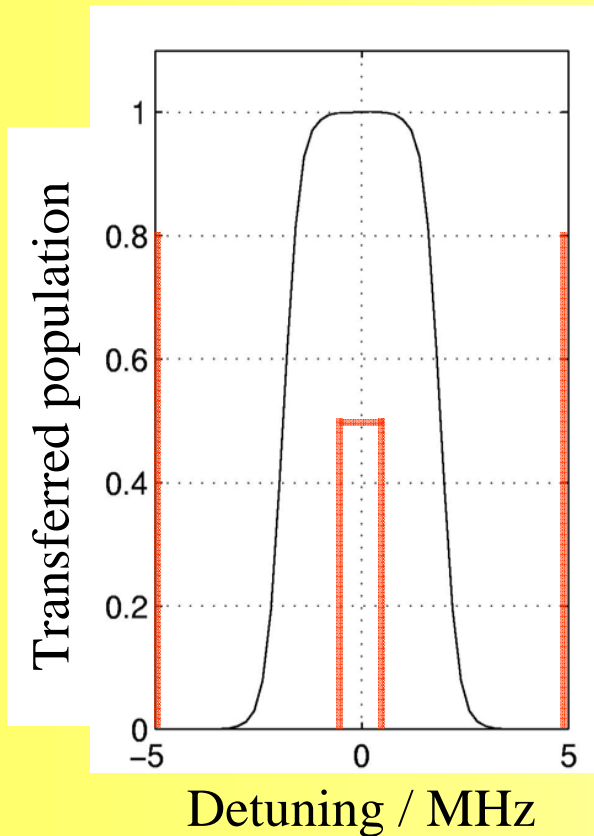
# Selecting strongly interacting ions



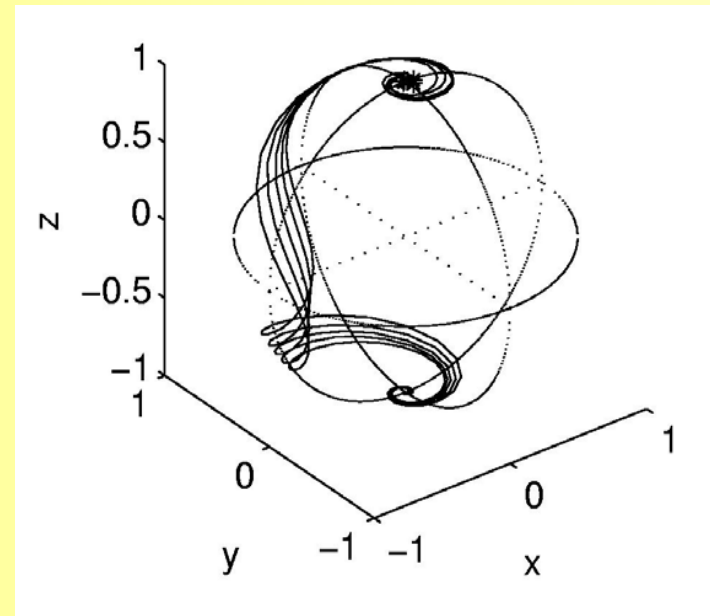
# Arbitrary single qubit operations

- Secant hyperbolic pulses compensate for the inhomogeneous broadening in the qubit for transitions going from one pole to the other pole on the Bloch sphere
- Thus they are not readily applicable to arbitrary single qubit operations

# Excitation with complex hyperbolic secant pulse



Evolution on the Bloch sphere



Further more, above a certain threshold intensity the operation is insensitive to different ions having different Rabi frequencies

# Pulse sequences for arbitrary single qubit operations

This work was carried out by Ingela Roos together with Klaus Mølmer

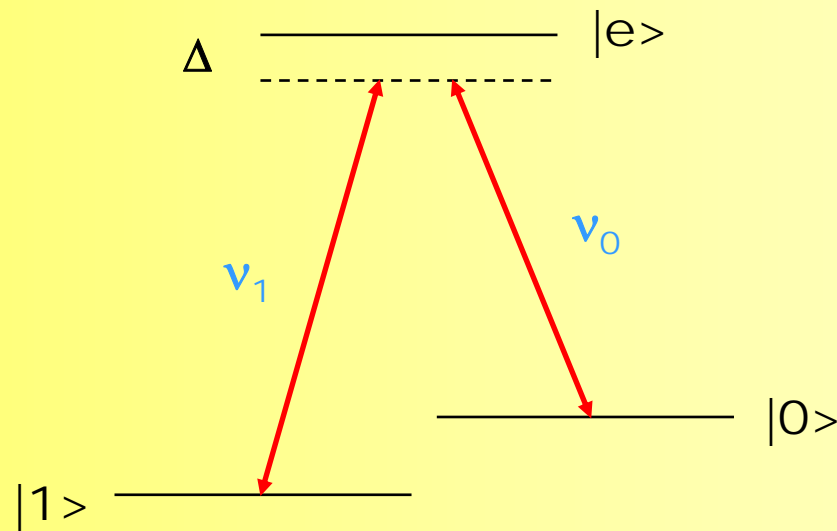
- “Robust quantum computing with composite pulse and coherent population trapping”, Phys Rev A **69**, 22321 (2004)



## Requirements

- Find pulses compensating for instances having different transition frequencies and different ion-field coupling strength

# Three-level system



$$\dot{c}_e = -i\Delta c_e + i\frac{\Omega_{R0}(t)}{2}e^{-i\varphi_0}c_0 + i\frac{\Omega_{R1}(t)}{2}e^{-i\varphi_1}c_1$$

# Three-level system

$$\Omega_{R0} = \Omega_{R1} \equiv \Omega_R$$

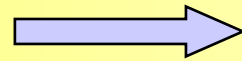
$$\longrightarrow \dot{c}_e = -i\Delta c_e + i \frac{\Omega_R(t)}{2} \underbrace{(e^{-i\varphi_0} c_0 + e^{-i\varphi_1} c_1)}_{=0}$$

= 0

$$\varphi_1 - \varphi_0 = \phi$$

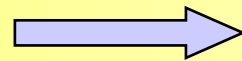
Coherent trapping

$$|d\rangle = \frac{1}{\sqrt{2}} (|0\rangle - \exp(-i\phi)|1\rangle)$$



$|d\rangle$  is a dark state

$$|b\rangle = \frac{1}{\sqrt{2}} (|0\rangle + \exp(-i\phi)|1\rangle)$$



$|b\rangle$  is a bright state

# Single qubit operations

unitary rotation about any axis on the equator of the Bloch sphere representing the qubit state basis

A two-colour sechyp pulse with phases  $\varphi_0$  and  $\varphi_1$  will drive the system into the excited state along route **A**.

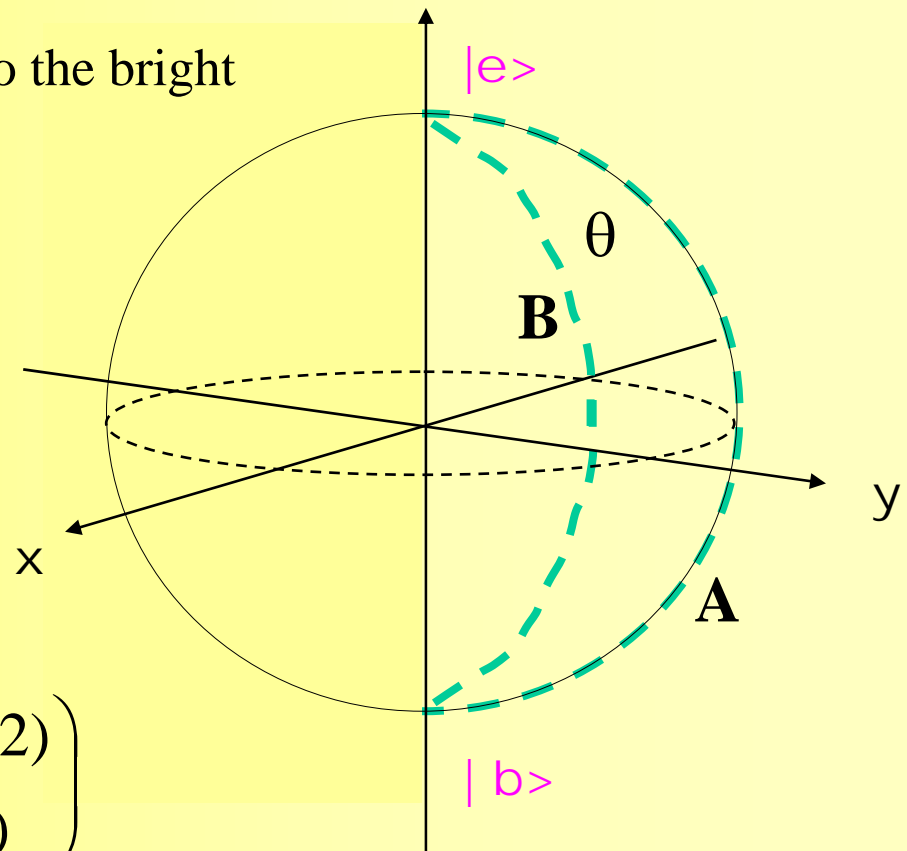
A new two-colour sechyp pulse with phases  $\varphi_0 + \pi + \theta$  and  $\varphi_1 + \pi + \theta$  will drive the system to the bright state along route **B**.

This corresponds to the operation

$$|b\rangle \rightarrow e^{i\theta} |b\rangle$$

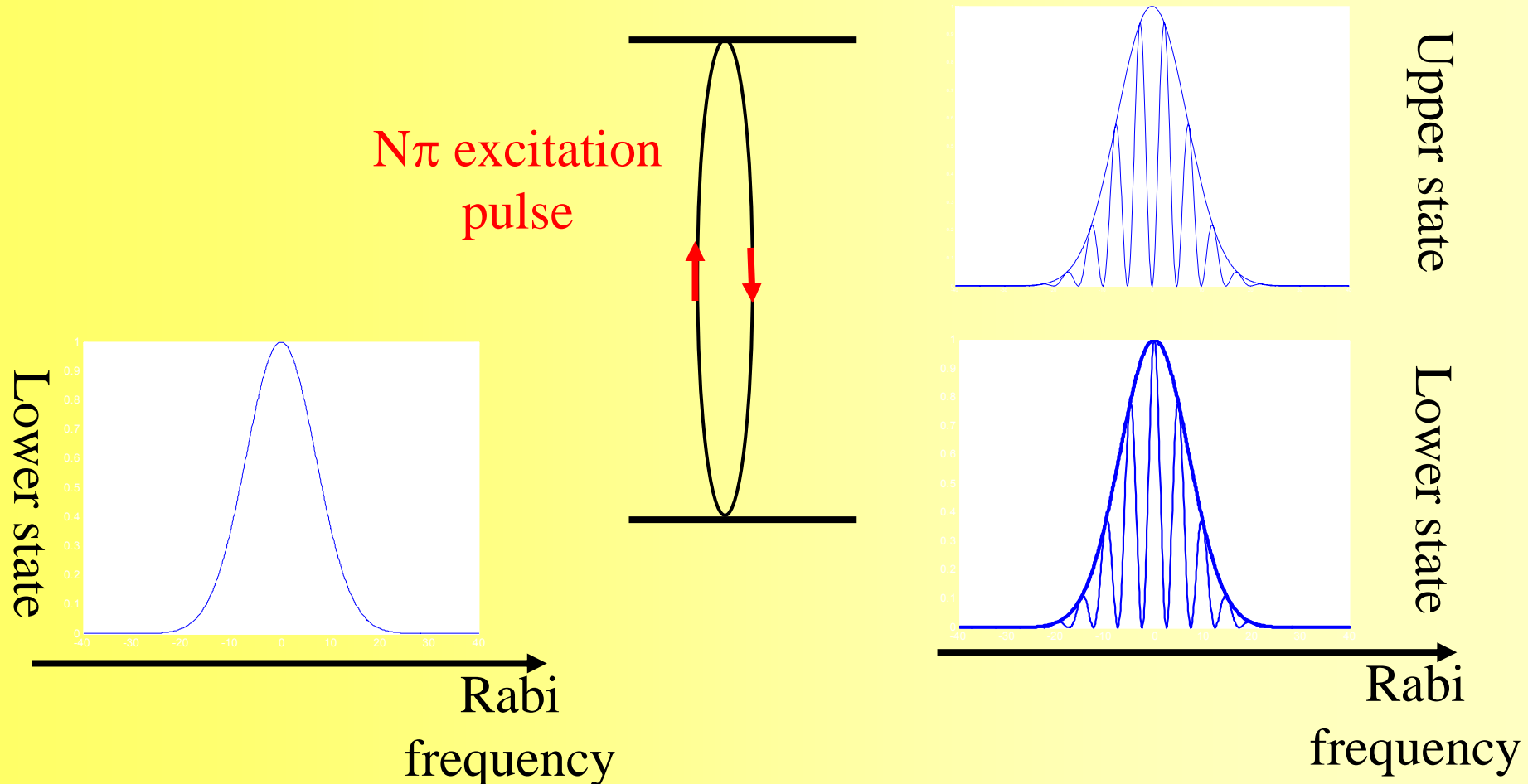
Which in the original  $|0\rangle, |1\rangle$  basis corresponds to the unitary operation

$$U = e^{i\theta/2} \begin{pmatrix} \cos(\theta/2) & ie^{i\phi} \sin(\theta/2) \\ ie^{-i\phi} \sin(\theta/2) & \cos(\theta/2) \end{pmatrix}$$



# Rabi frequency selection

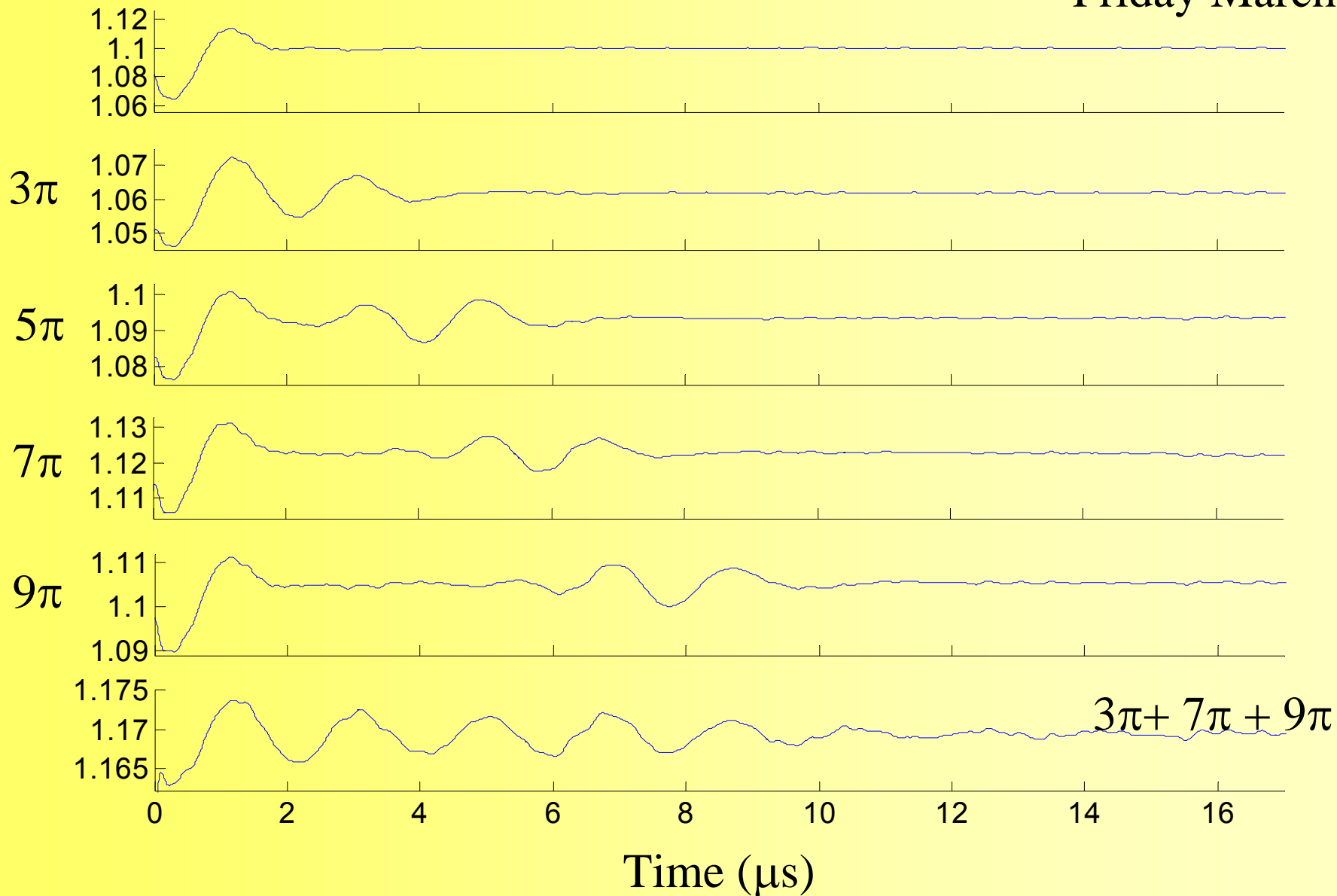
Ions in different parts of the beam experience different Rabi frequencies





# Rabi frequency distillation

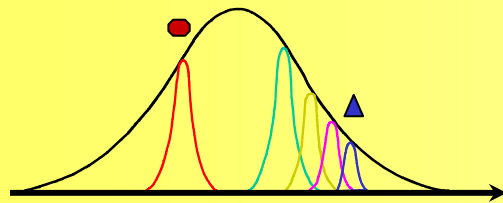
Friday March 9



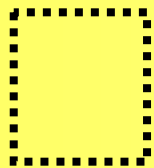
# Requirements for quantum computing

- Coherent two-level systems acting as qubits
- Possibility to manipulate the qubits individually (single qubit operations)
- Coupling between any two qubits (two-bit gates)
- Possibility for reliable read-out of the individual qubits
- **Scalability**

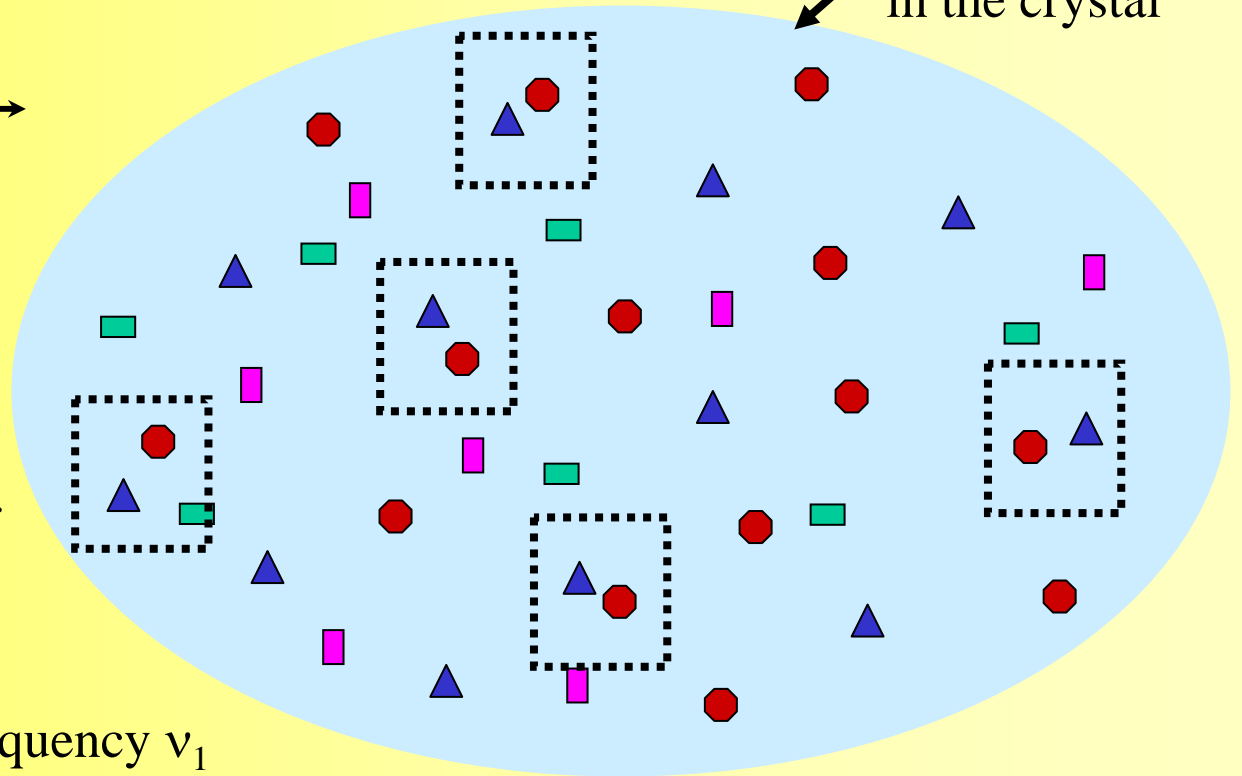
# Qubit distillation



Arbitrary volume  
in the crystal



Ions interacting  
strongly enough  
for mutual control.  
Potential QC



● ion absorbing at frequency  $\nu_1$

▲ ion absorbing at frequency  $\nu_2$

# Dipole-dipole interaction

$\delta\nu$  – frequency shift due to the dipole-dipole interaction

$$\delta\nu \propto \frac{(\Delta\mu)^2}{r^3}$$

$\Delta\mu$  - difference in dipole moment between ground and excited state

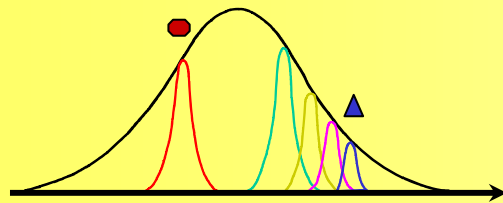
$r$  – distance between interacting ions

Fraction of controllable ions in a qubit ( $\eta$ ) scale as  $\eta \propto N(\Delta\mu)^2$

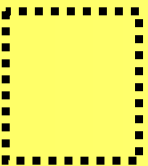
$N$  = dopant concentration


Higher dopant concentration, or larger difference in dipole moments between ground and excited state would lead to increased probability to find ions that interact


# Scaling

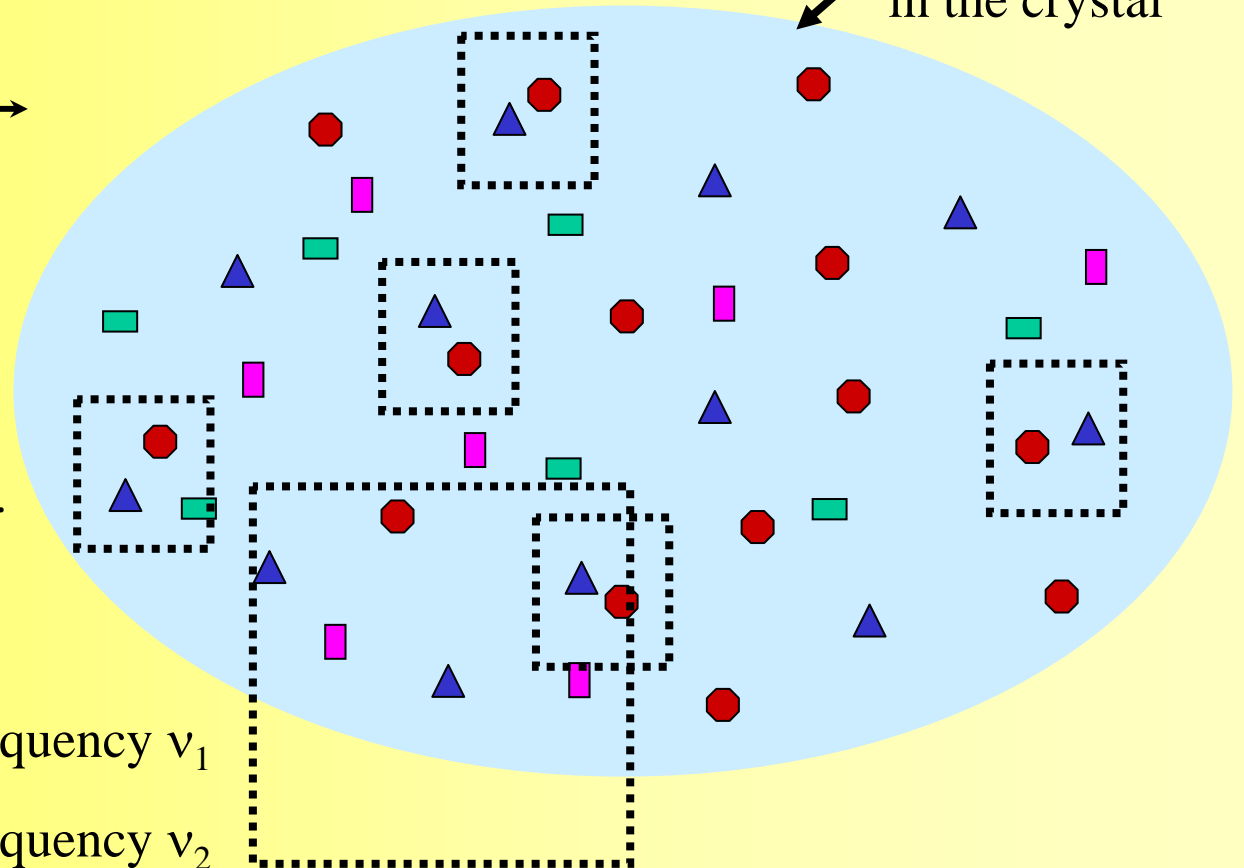


Arbitrary volume  
in the crystal

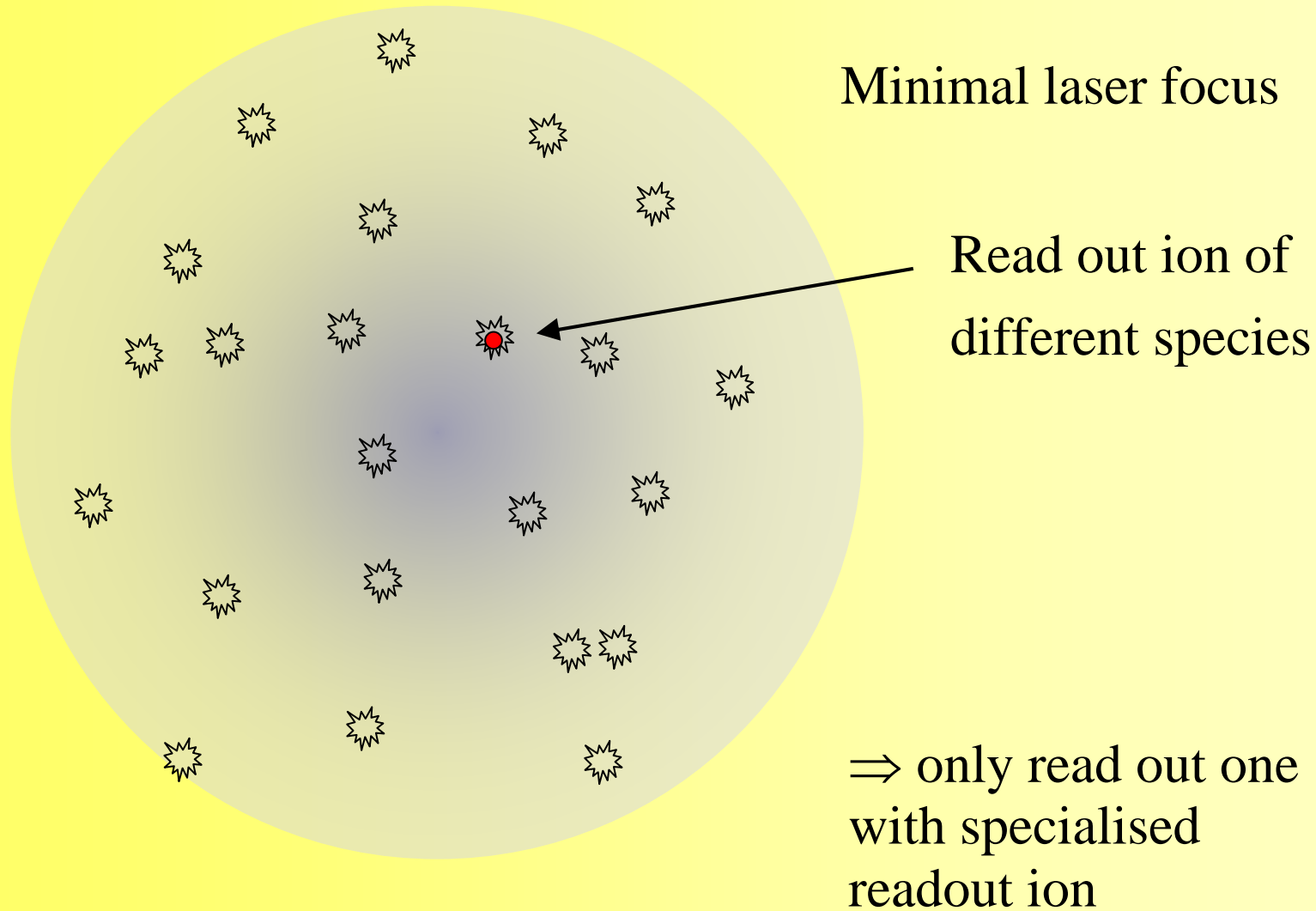
 Ions interacting  
strongly enough  
for mutual control.  
Potential QC

 ion absorbing at frequency  $\nu_1$

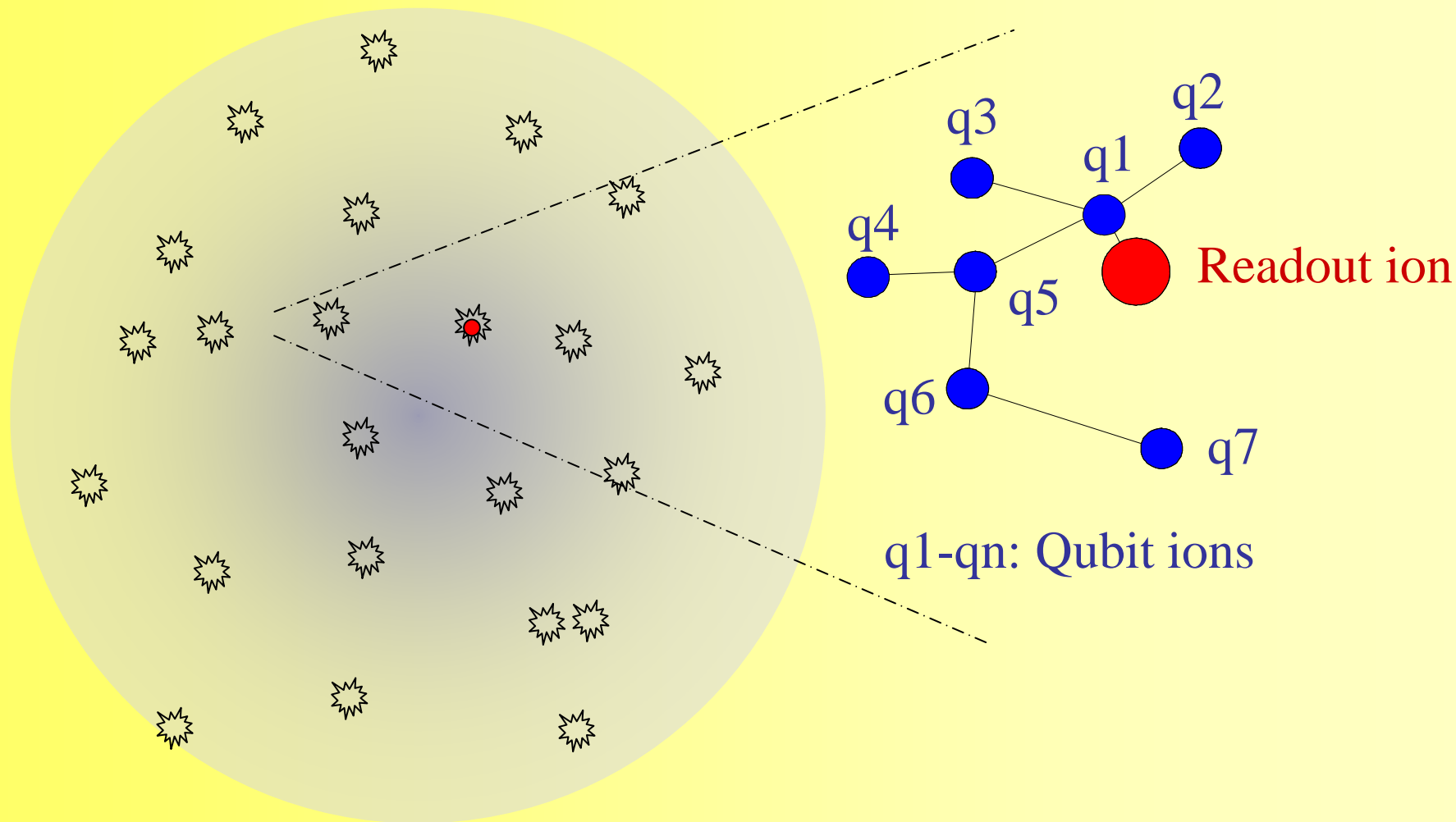
 ion absorbing at frequency  $\nu_2$



# Readout: Single instance with search approach



# Single instance with search approach



# Summary

## Selected results

- Qubits are initiated in well defined states
- Qubit state-to-state transfer efficiency is 97.5% , simulations predict 98.5%
- Qubit distillation >90%
- Scalable schemes have been developed



# Outlook

- Our dye laser system has been frequency stabilised to a few kHz linewidth to carry out high fidelity:
  - Single qubit operations
  - Two-qubit gate operations
- Potential read out ion candidates are investigated



Atia Amari



Lars Rippe



Andreas Walther



Yan Ying

## Former members



Brian  
Julsgaard



Mattias  
Nilsson



Nicklas  
Ohlsson



Ingela  
Roos