

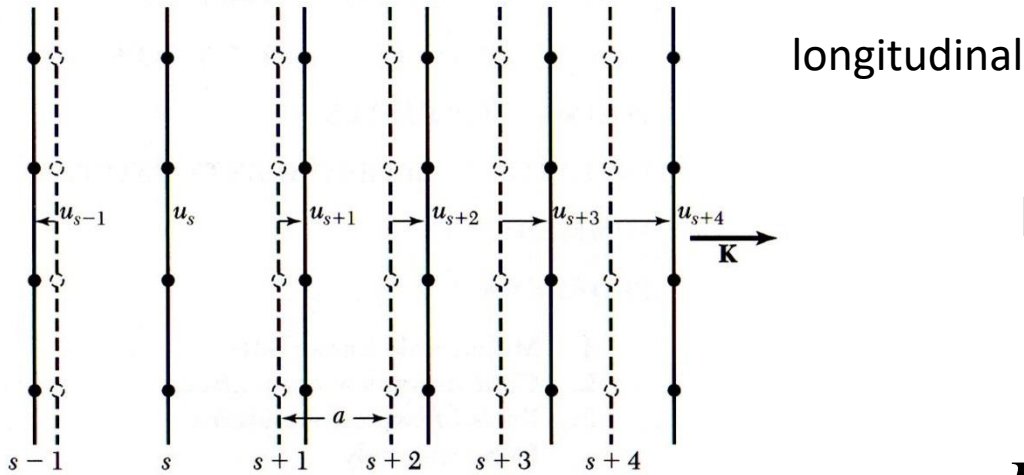
Solid state physics for Nano



Lecture 4: Phonons

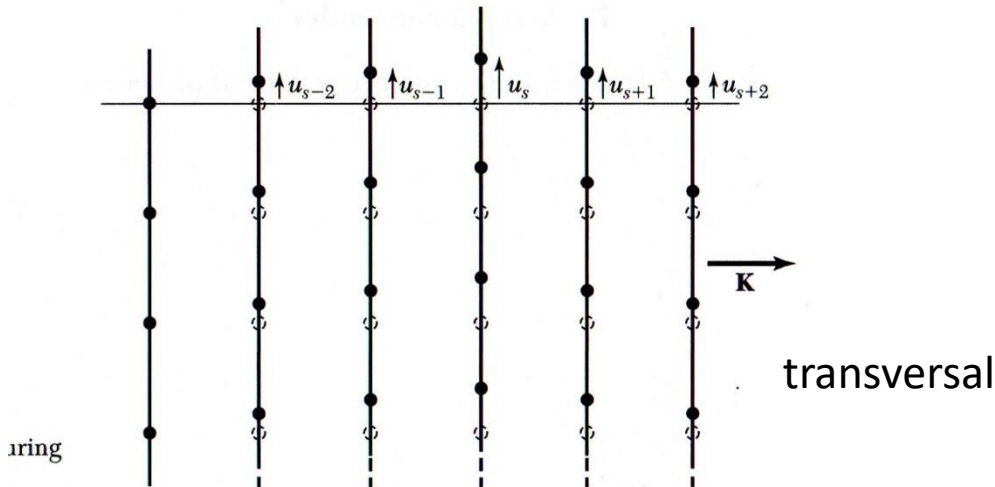
Prof. Dr. U. Pietsch

Crystal lattice vibration



Hook's law

$$F_n = \sum_p C_p (u_{n+p} - u_n)$$



F_n – Force on plane n ;
 C_p – force constant between planes
of distance p

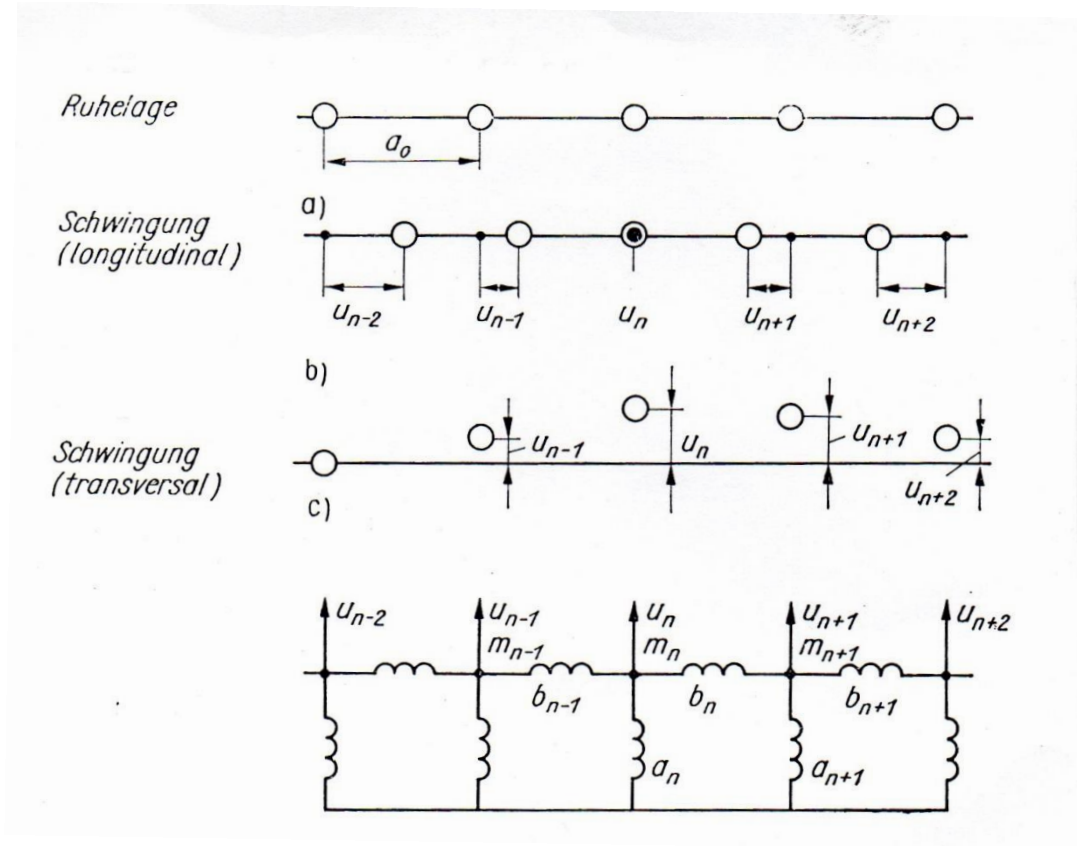
lattice vibration of 1D monoatomic chain

Equation of motion

$$M \frac{d^2 u_n}{dt^2} = \sum_{p=-\infty}^{+\infty} C_p (u_{n+p} - u_n)$$

Ansatz:

$$u_{n+p} = u_0 e^{i[(n+p)\vec{q}\vec{a} - \omega t]}$$



a – lattice spacing; q – wave number, \vec{q} - wave vector

Solution of 1D wave equation

$$\rightarrow(*) \quad M \frac{d^2 u_n}{dt^2} = -\omega^2 M u_0 e^{i[(nqa - \omega t)]} = \sum_p [e^{i[(n+p)qa]} - e^{i[nqa]}] e^{-i\omega t}$$

$$\omega^2 M = -\sum_p C_p (e^{ipqa} - 1)$$

$$C_{-p} = -C_p \quad \omega^2 M = -\sum_{p>0} [C_p (e^{ipqa} - 1) + C_p (e^{-ipqa} - 1)]$$

$$\omega^2 M = -\sum_{p>0} C_p (e^{ipqa} + e^{-ipqa} - 2) = -\sum_{p>0} C_p (2 \cos(pqa) - 2)$$

$$\text{Dispersion relation} \quad \omega^2 = \frac{2}{M} \sum_{p>0} C_p (1 - \cos(pqa))$$

$\omega(q)$ – dispersion – nonlinear function

Only nearest neighbours , p=1

$$\omega^2 = \frac{2C_1}{M} (1 - \cos(pa)) = \frac{2C_1}{M} \left(2 \sin^2 \frac{qa}{2}\right)$$

$$\omega^2 = \frac{4C_1}{M} \sin^2\left(\frac{qa}{2}\right)$$

Small q

$$\omega^2 = \frac{2C}{M} (1 - \cos(pqa)) = \frac{2C}{M} \left(1 - 1 + \frac{1}{2} q^2 p^2 a^2 \dots\right)$$

$$\omega^2 = \frac{C_1}{M} q^2 a^2 \quad \omega \sim q$$

1st Brillouin zone contains all information : $\omega(q)$ is periodic in q^*a

$$\frac{u_{n+1}}{u_n} = \frac{e^{i[(n+1)qa - \omega t]}}{e^{i[nqa - \omega t]}} = e^{i(qa - \omega t)}$$

Only in range $|qa| < 2\pi$ independent values of q

$$-\frac{\pi}{a} \leq q \leq \frac{\pi}{a}$$

Range of 1st Brillouin zone

$$\omega^2 = \frac{4C_1}{M} \sin^2\left(\frac{qa}{2}\right)$$

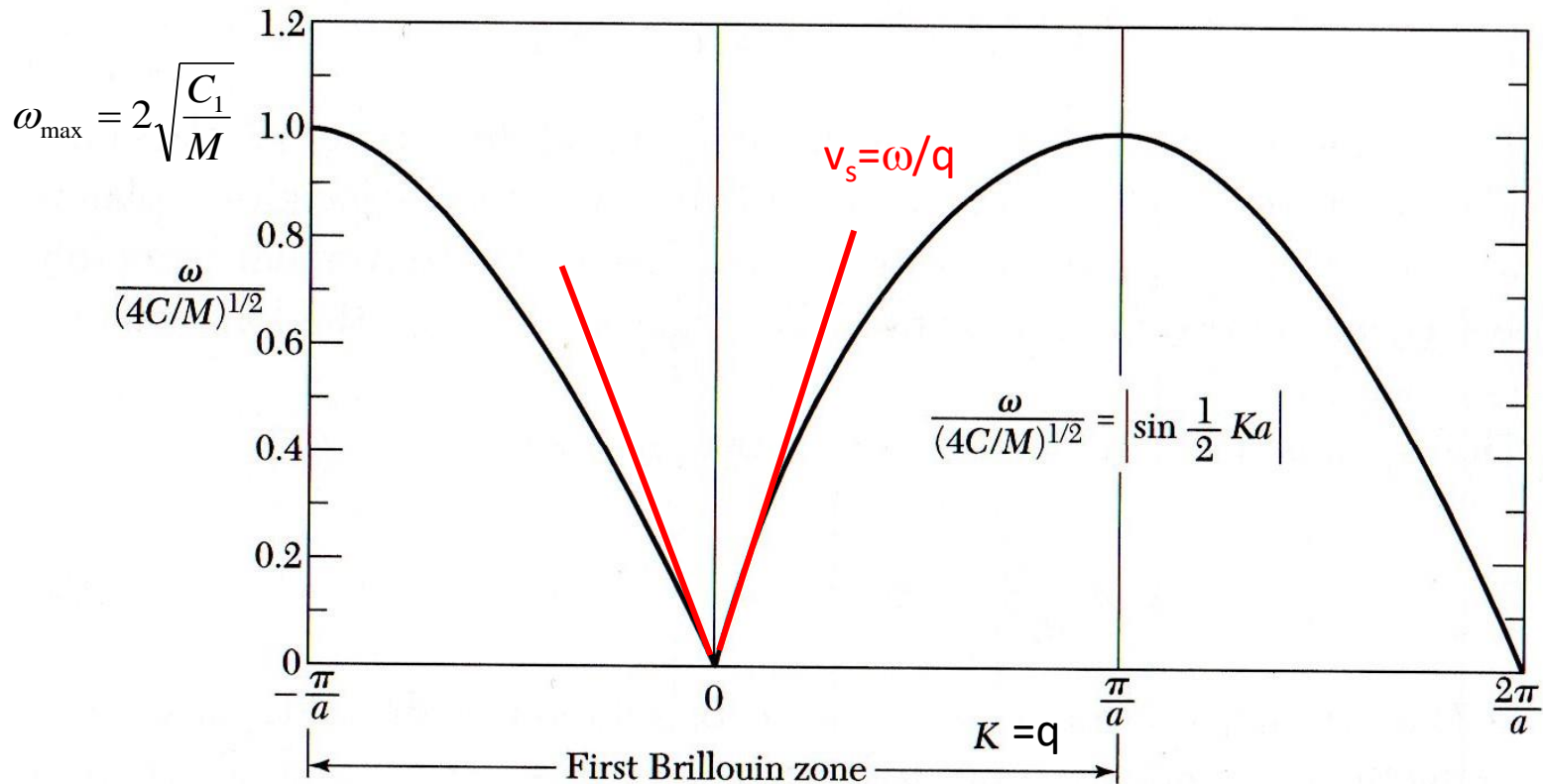


Figure 4 Plot of ω versus K . The region of $K \ll 1/a$ or $\lambda \gg a$ corresponds to the continuum approximation; here ω is directly proportional to K .

$$q_{\max} = \pm \frac{\pi}{a} \approx 10^{10} \text{ m}^{-1}$$

$$\lambda_{\min} = \frac{2\pi}{q_{\max}} = 2a$$

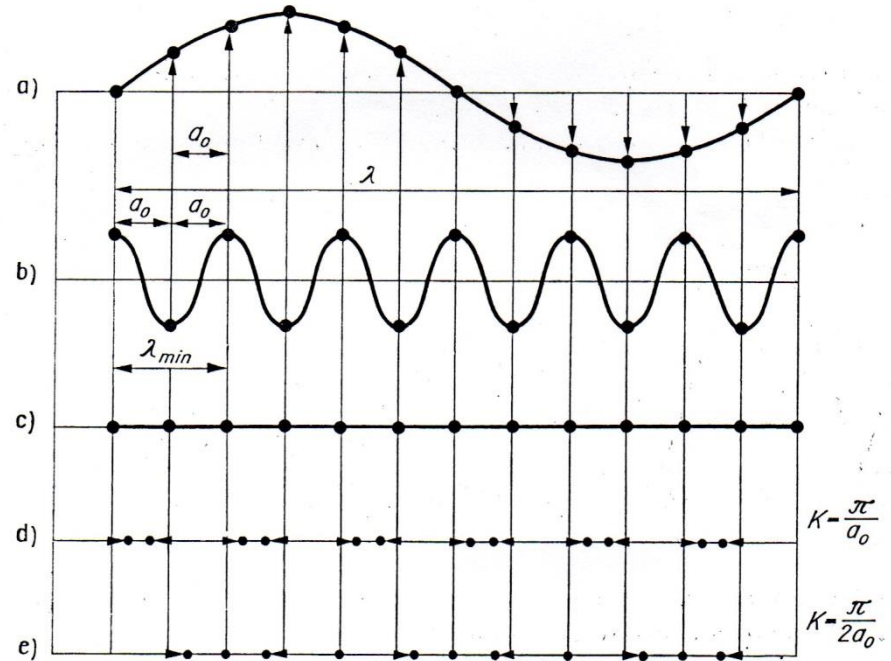


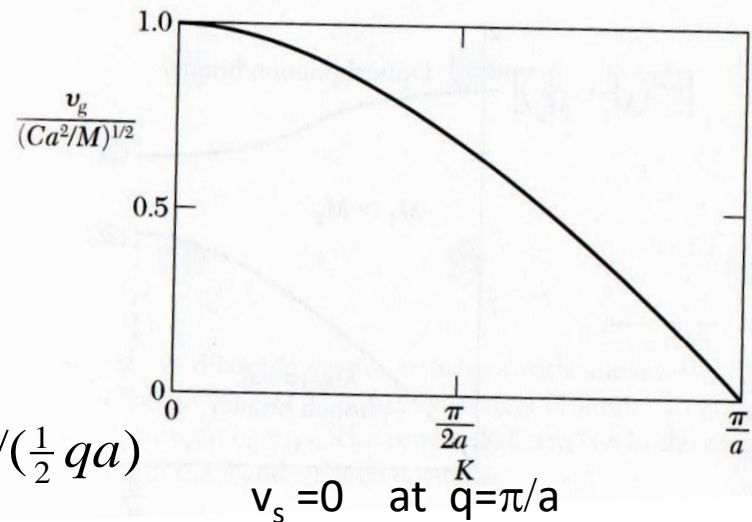
Bild 2.14. Schwingungen linearer Atomketten

Group velocity = signal velocity of a wave packet

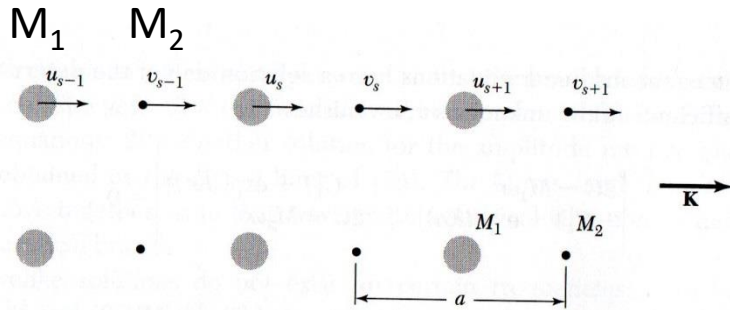
$$v_g = \frac{d\omega}{dq} \quad v_g = \sqrt{\frac{Ca^2}{m}} \cos\left(\frac{1}{2} qa\right)$$

phase velocity > group velocity

$$v_p = \frac{\omega}{q} \quad v_p = \sqrt{\frac{Ca^2}{m}} \sin\left(\frac{1}{2} qa\right) / \left(\frac{1}{2} qa\right)$$



lattice vibration of 1D bi-atomic chain



$$M_1 \frac{d^2 u_s}{dt^2} = C(v_s + v_{s-1} - 2u_s)$$

$$M_2 \frac{d^2 v_s}{dt^2} = C(u_{s+1} + u_s - 2v_s)$$

$$u_s = u_0 e^{i(sqa - \omega t)} \quad v_s = v_0 e^{i(sqa - \omega t)}$$

$$-\omega^2 M_1 u = Cv [1 + e^{iqa}] - 2Cu$$

$$-\omega^2 M_2 v = Cu [e^{iqa} + 1] - 2Cv$$

$$\begin{vmatrix} 2C - M_1 \omega^2 & -C[1 + e^{iqa}] \\ -C[1 + e^{iqa}] & 2C - M_2 \omega^2 \end{vmatrix} = 0$$

Quadratic equation in ω^2

$$M_1 M_2 \omega^4 - 2C(M_1 + M_2) \omega^2 + 2C^2(1 - \cos qa) = 0$$

$$\omega^2 = C\left(\frac{1}{M_1} + \frac{1}{M_2}\right) \pm C\sqrt{\left(\frac{1}{M_1} + \frac{1}{M_2}\right)^2 - \frac{4\sin^2 qa}{M_1 M_2}}$$

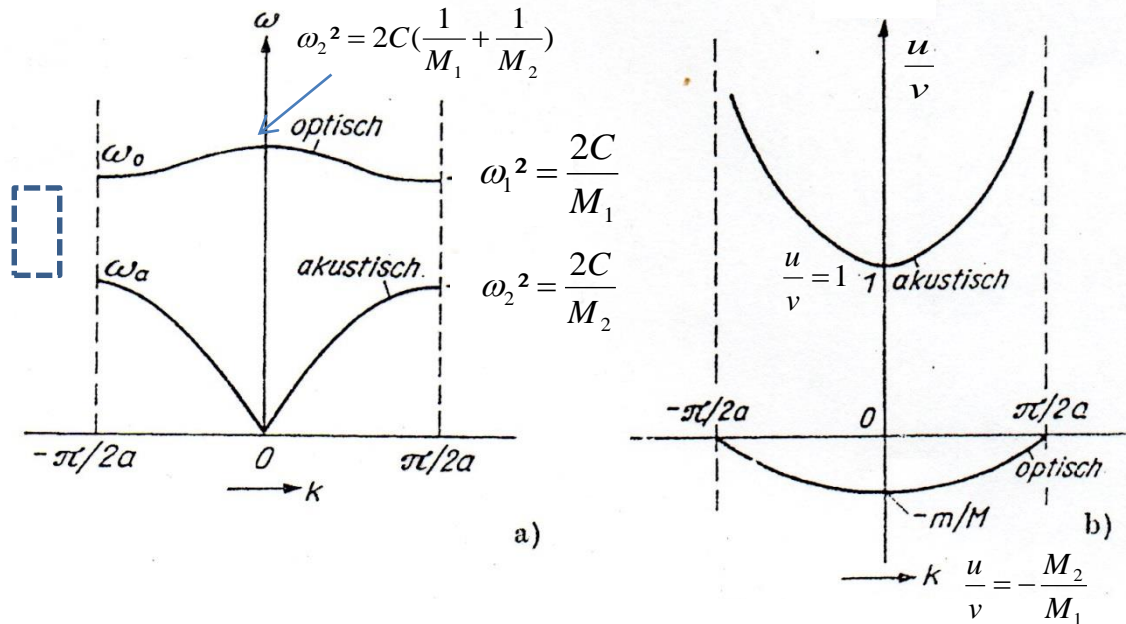
Acoustic branch

For $q=0$ $\omega_1^2 = 0$ For $\sin(qa) \approx qa$ $\omega_1^2 = \frac{\frac{1}{2}C}{M_1 + M_2} (qa)^2$

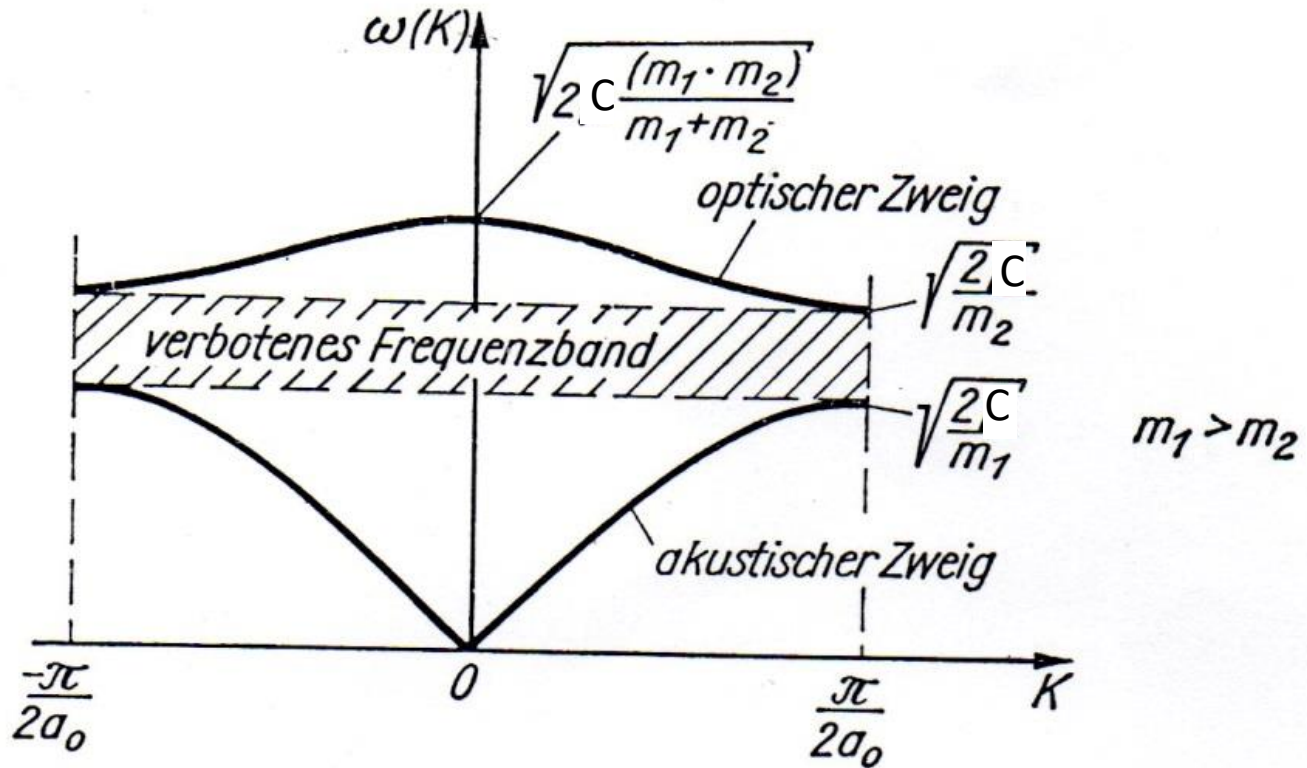
Optical branch

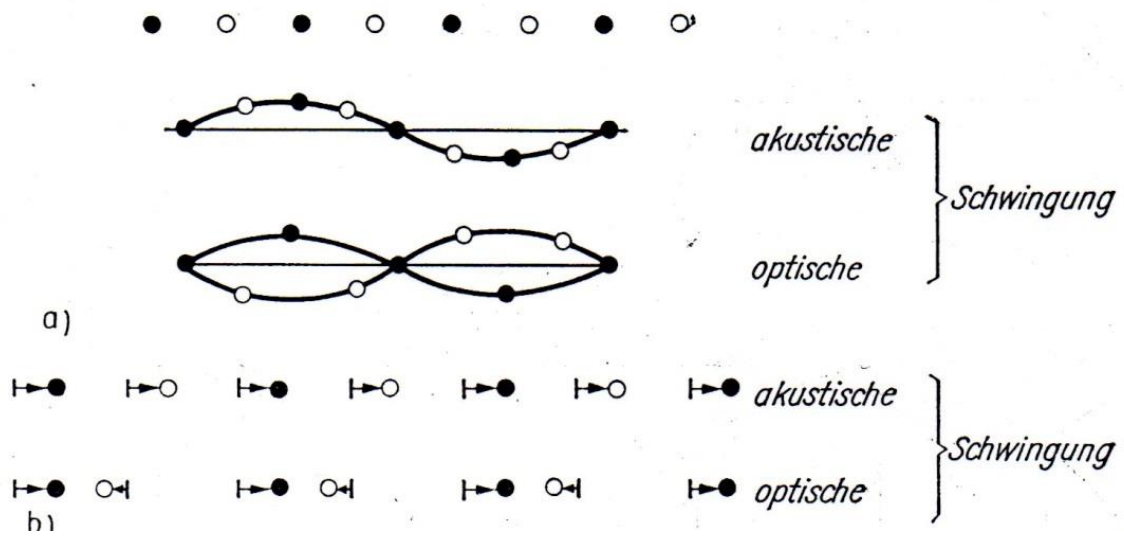
$$\omega_2^2 = 2C\left(\frac{1}{M_1} + \frac{1}{M_2}\right)$$

Optical frequency gap

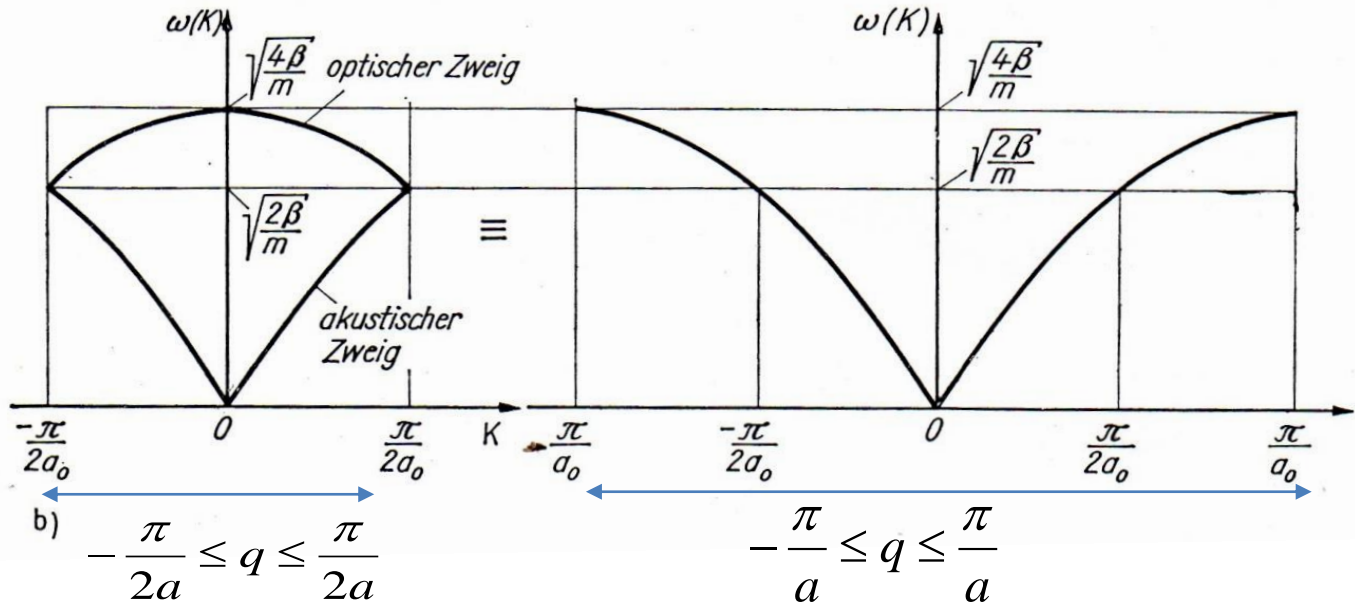


	$q = 0$	$q = \pm \frac{\pi}{2a}$
<i>acoustic_branch</i>	$\omega = 0; u = v$	$\omega = \sqrt{\frac{2C}{M_1}}; v = 0$; $u_{n-1} = -u_{n+1}$
<i>optical_branch</i>	$\omega = \sqrt{2C\left(\frac{1}{M_1} + \frac{1}{M_2}\right)}, \frac{u}{v} = -\frac{M_2}{M_1}$	$\omega = \sqrt{\frac{2C}{M_2}}; u = 0; v_{n-1} = -v_{n+1}$





Equal masses : $M_1=M_2$



Phonons in 3D

- 3D crystal of size Pa_1, Pa_2, Pa_3 may contain S different atoms per unit cell
- Each atoms has $3SP^3=3SN$ degrees of freedom for displacements
- Each displacement results in 1D equations of motion
- \rightarrow $3SN$ differential equation with $3SN$ solutions
- Direction of wave propagation given by vector \mathbf{K}
- Displacement $u = u_0 \exp(i(\mathbf{K}'\mathbf{a}_n - \omega t))$ using

$$\vec{K}' = \vec{K} + \vec{G} \quad \vec{G} = hb_1 + kb_2 + lb_3$$

$$\vec{K}'\vec{a}_n = \vec{K}\vec{a}_n + (hb_1 + kb_2 + lb_3)\vec{a}_n = \vec{K}\vec{a}_n + m2\pi \quad \mathbf{K}' \text{ periodic in 3D}$$

- Valid values in 1st BZ
$$-\pi \leq \vec{K}\vec{a}_n \leq \pi$$
- Number of solutions is $3S$
- $3S$ equation give $3S$ branches $\omega_j(\mathbf{K}), j=1,2,3$, corresponds to 1 long + 2 trans modes
- In total 3 acoustic branches and $3(S-1)$ optical branches

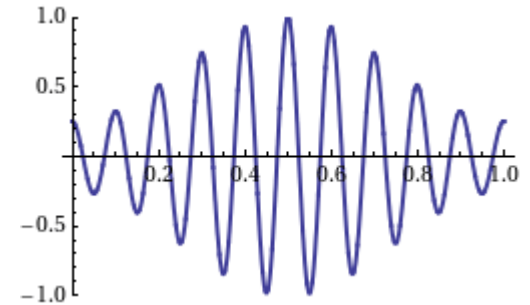
Group and phase velocities

Group velocity = signal velocity

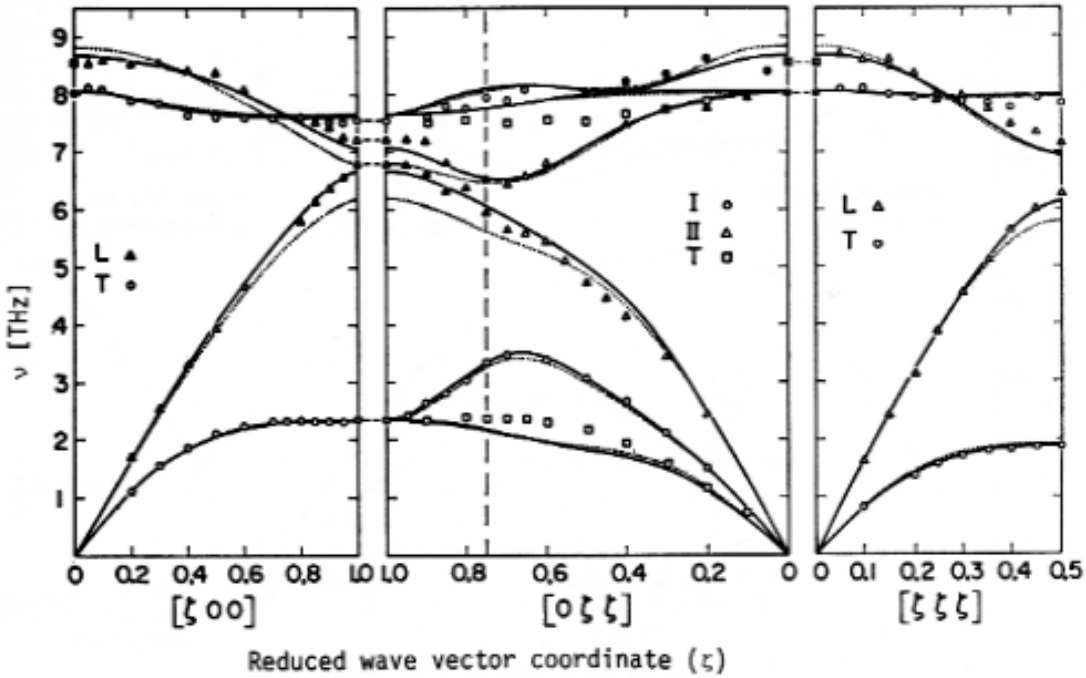
$$v_{gr} = \frac{d\omega}{dq} = v_0 \left(\cos \frac{qa}{2} \right)$$

Phase velocity

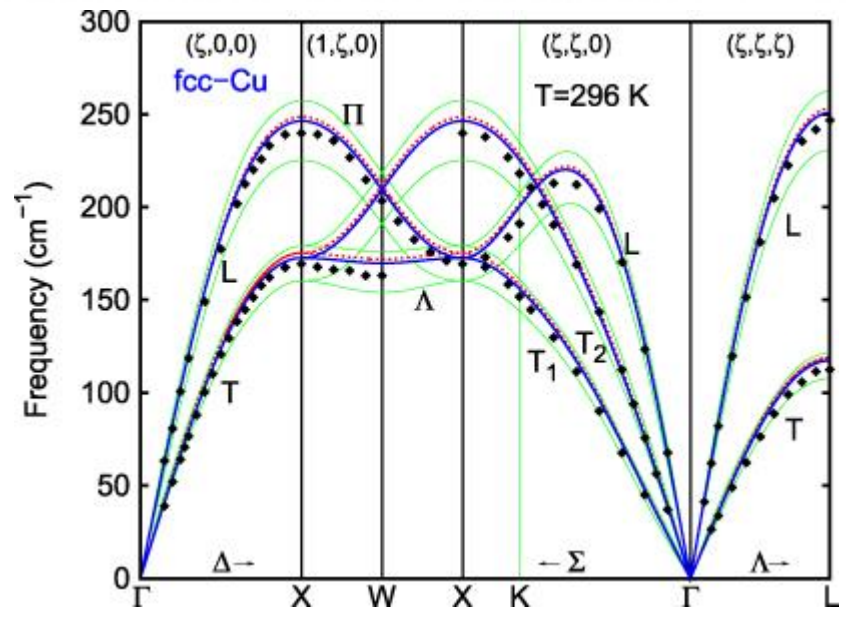
$$v_{ph} = \frac{\omega}{q} = v_0 \left| \frac{\sin(qa/2)}{qa/2} \right|$$



This shows a wave with the group velocity and phase velocity going in different directions. The group velocity is positive, while the phase velocity is negative.



Phonon dispersion of GaAs



***Ab initio* phonon dispersions of transition and noble metals: effects of the exchange and correlation functional**
Andrea Dal Corso

Phonon dispersion of Cu